Structural phase transition in the two-dimensional triangular lattice antiferromagnet RbFe(MoO$_4$)$_2$

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X-ray diffraction data were collected on powder samples of RbFe(MoO$_4$)$_2$, grown by a flux method described in Ref. 3. Greenish, transparent thin crystals with typical dimensions of $2 \times 2 \times 0.2$ mm$^3$ were obtained. X-ray diffraction data were collected on powder samples of RbFe(MoO$_4$)$_2$, grown by a flux method described in Ref. 3. Greenish, transparent thin crystals with typical dimensions of $2 \times 2 \times 0.2$ mm$^3$ were obtained.

The behavior of the magnetization and ESR spectra correspond roughly to that of a planar antiferromagnet with the "triangular" magnetic structure. Nevertheless, a weak interplane interaction gives rise to additional phase transitions and spin-wave modes. Moreover, the magnetic structure appears to be incommensurate in the $c$ direction (perpendicular to magnetic planes). The nature of the incommensurate modulation of the magnetic structure is still unknown.

Several members of the aforementioned family of trigonal molybdates and tungstates undergo structural phase transitions with temperature. As far as we know, the crystal structure of RbFe(MoO$_4$)$_2$ was investigated only at room temperature.

The knowledge of the distinct crystal structure is of importance for the explanation of the incommensurate magnetic structure, because the appropriate phase transitions are symmetry sensitive (see, e.g., Ref. 8). In this work, we study the temperature-dependent peculiarities of the structure of RbFe(MoO$_4$)$_2$ using x-ray diffraction, electron spin resonance (ESR), and Raman scattering.

II. EXPERIMENT

Single crystals of RbFe(MoO$_4$)$_2$ were grown by a flux method described in Ref. 3. Greenish, transparent thin crystals with typical dimensions of $2 \times 2 \times 0.2$ mm$^3$ were obtained.
prepared from single crystals, at the temperatures 300, 250, 200, 180, 150, and 100 K. CuKα radiation was monochromatized by a curved germanium crystal. The indexation of the powder x-ray patterns was made using the program WINXPOW-1999. ESR measurements in the temperature range \(4.2 < T < 300\) K were performed with a Bruker ELEXSYS E500CW spectrometer at X-band and Q-band frequencies \(9.5\) and \(34\) GHz, respectively. The spectrometer was equipped with an Oxford Instruments continuous He flow cryostat. The magnetic susceptibility at the temperatures between 1.8 and 300 K was measured with a SQUID magnetometer.

Raman scattering was excited by the 514-nm line of an Ar-ion laser. Raman spectra at room temperature were measured in the Institute of Spectroscopy, employing a homemade triple spectrograph and a multichannel system consisting of an image-intensifier tube with a multichannel plate and a vidicon. The temperature-dependent Raman spectra were taken in the University of Groningen using a SPEX 14018 spectrometer and a Burle C13-34 photomultiplier. The typical laser power in these experiments was 20 mW. A single crystal of RbFe(MoO₄)₂ was attached to a copper sample holder by a silver paste and was placed into a closed-cycle He cryostat. The temperature of the holder near the sample location was measured with the aid of a silicon diode and could be varied between 33 and 300 K. Most of the samples strongly depolarize the light, so it was difficult to perform polarization measurements. Partially polarized Raman scattering in \(z(\pm y)\) and \(z(x)\) configurations has been measured from a selected sample at room temperature.

### III. RESULTS

#### A. X-ray powder diffraction

Powder x-ray diffraction patterns taken at 300 and 100 K are shown in Fig. 2(a). These patterns are similar one to another, the last one demonstrates some shift to larger angle values. At 180 K and, to a less extent, at 150 K several reflections became asymmetric. The temperature dependence of the unit-cell parameters \(a\) and \(c\) display a peculiarity near 200 K pointing to a phase transition [see Fig. 2(b)]. Unfortunately, the accuracy of our x-ray experiment at low temperatures did not allow to make a choice between the close space groups \(P\overline{3}m1\) and \(P\overline{3}c1\).

#### B. Electron spin resonance

At temperatures higher than the Néel temperature, the ESR spectrum of the RbFe(MoO₄)₂ crystal consists of a single exchange-narrowed Lorentzian line with a resonance field corresponding to a \(g\) factor near \(g \approx 2\). Figure 3 illustrates the results obtained at the microwave frequency 9.48 GHz (X band). The upper frame shows the temperature dependence of the resonance linewidth \(\Delta H\) [here \(\Delta H\) is the half-width at half maximum (HWHM)]. Since the linewidth is of the same order of magnitude as the resonance field (shown in the lower frame), the spectra were evaluated taking into account both circular components of the resonance, as described, e.g., in Ref. 9. The linewidth exhibits a pronounced anisotropy with respect to the crystallographic \(c\) axis. For the external magnetic field \(H\) applied parallel to the \(c\) axis \((\theta = 0\,^\circ)\), the width is two times larger (1.5 kOe at room temperature) than for the magnetic field applied within the basal plane \((\theta = 90\,^\circ)\), where it remains unchanged on rotation about the \(c\) axis. With decreasing temperature, the linewidth gradually increases and finally, near \(T_N\), demonstrates a rapid growth typical for a critical behavior. Focusing on the phase transition near 200 K, where the lattice
constant changes, we observe a distinct step of about 4% in the linewidth at $T_c = 190$ K. At the same temperature, the resonance field $H_{\text{res}}$, which remained nearly unchanged above 50 K, exhibits a shift of approximately 1%. The inverse ESR intensity $1/I_{\text{ESR}}$ agrees well with the Curie-Weiss law found for the inverse susceptibility $1/x$ from SQUID measurements, as depicted in the inset of Fig. 3. Both show a slight kink near $T_c$, as well.

The critical line broadening on approaching $T_N$ indicates an increasing importance of short-range order effects below 50 K. These effects give rise to fluctuating internal fields, which are independent of the external field, and which shift the resonance field. This shift is added to the paramagnetic shift due to the single ion anisotropy, and will be discussed below. To separate both contributions to the resonance shift, we performed additional ESR measurements at a higher microwave frequency of 34 GHz ($Q$ band). Figure 4 shows the temperature dependences of the resonance field (lower frame) and the linewidth (upper frame) for both main orientations of the magnetic field at 34 GHz. The linewidth data agree well with those obtained at 9.48 GHz (see inset). The anisotropy of the resonance field increases with decreasing temperature. The clear decrease of the resonance field observed at 34 GHz for the magnetic field applied in the basal plane was not detected at 9.48 GHz, but just the opposite behavior was observed resulting from the dominance of the short-range order in low external fields. However, at 34 GHz the paramagnetic shift governs the resonance position due to the higher external field.

C. Group-theoretical analysis of vibrational modes

RbFe(MoO$_4$)$_2$ contains one formula unit ($N = 12$ atoms) in a primitive cell of the trigonal space group $P\overline{3}m1$. Thus, there are $3N - 3 = 33$ optical vibrations. To determine their symmetries, we perform a nuclear-site group analysis. Table I displays the crystallographic positions occupied by different atoms and the contributions from these atoms to the vibrational modes of the crystal. The atomic position is indicated according to the Wyckoff notation. The number of equivalent sites is given by a number before the letter notation.

<table>
<thead>
<tr>
<th>Atom</th>
<th>Local symm.</th>
<th>Vibrational modes</th>
</tr>
</thead>
<tbody>
<tr>
<td>Rb</td>
<td>$1b$</td>
<td>$D_{3d}$</td>
</tr>
<tr>
<td>Fe</td>
<td>$1a$</td>
<td>$D_{3d}$</td>
</tr>
<tr>
<td>Mo</td>
<td>$2d$</td>
<td>$C_{3v}$</td>
</tr>
<tr>
<td>O1</td>
<td>$2d$</td>
<td>$C_{3v}$</td>
</tr>
<tr>
<td>O2</td>
<td>$6i$</td>
<td>$C_6$</td>
</tr>
</tbody>
</table>

$A_{2u} + E_u$

$A_{2u} + E_u$

$A_{1g} + A_{2u} + E_u$

$A_{1g} + A_{2u} + E_u$

$A_{1g} + A_{2u} + 3E_u + 3E_u$


FIG. 3. Temperature dependence of the 9.48 GHz ESR linewidth (upper frame) and the resonance field (lower frame) in RbFe(MoO$_4$)$_2$ for the magnetic field applied parallel ($\theta = 0^\circ$) and perpendicular ($\theta = 90^\circ$) to the crystallographic $c$ axis. Inset: temperature dependence of the inverse ESR intensity $1/I_{\text{ESR}}$ compared to the inverse static susceptibility $1/x$.

FIG. 4. Temperature dependence of the 34 GHz ESR linewidth (upper frame) and the resonance field (lower frame) in RbFe(MoO$_4$)$_2$ for the magnetic field applied parallel ($\theta = 0^\circ$) and perpendicular ($\theta = 90^\circ$) to the crystallographic $c$ axis. The solid line is obtained from the fit described in the discussion. Inset: Angular dependence of the linewidth at 150 K for both $X$- and $Q$-band frequencies.

TABLE I. Crystallographic positions occupied by different atoms and corresponding contributions to the vibrational modes of RbFe(MoO$_4$)$_2$, $P\overline{3}m1$ space group.
tion. The local symmetry of an atomic position is indicated by the point-group symbol (in Schönflies notation). We note that there are two nonequivalent oxygen atoms O1 and O2. O1 atoms reside directly above (or below), Mo atoms in the MoO4 tetrahedra, whereas O2 atoms are in the basal plane. Summarizing the contributions from all the atoms and subtracting the acoustic modes \( (A_{2g}+E_{u}) \), we obtain the following optical vibrational modes of RbFe(MoO4)2:

\[
\Gamma^{opt}(P\bar{3}m1) = 4A_{1g}(xx,yy,zz) + A_{2g} + 5E_{g}(xz,yz,xy) + A_{1u} + 5A_{2u}(E|x|z) + 6E_{u}(E|x,E|y). \tag{1}
\]

There are 9 Raman \( (4A_{1g} + 5E_{g}) \) and 11 infrared \( (5A_{2u} + 6E_{u}) \) modes. Note that the space group of RbFe(MoO4)2 contains the inversion operation and therefore Raman and infrared modes are well separated. Even \( (g) \) vibrations are Raman active, whereas the odd \( (u) \) ones are infrared active. \( A_{2g} \) and \( A_{1u} \) modes are silent.

Numerous studies of different molybdates have shown that a part of the crystal vibrations may be identified with the internal motions of the “free” \( (\text{MoO}_4)^{2-} \) tetrahedral ion (which has \( T_{d} \) symmetry), \( v_{1}(A_{1}) \), \( v_{2}(E) \), \( v_{3}(F_{2}) \), and \( v_{4}(F_{2}) \).\(^{11,12} \) As usual, \( A, E, \) and \( F \) letters refer to nondegenerate, doubly and triply degenerate vibrations, respectively. To show the origin of the crystal lattice vibrations, we construct a correlation scheme for RbFe(MoO4)2 (see Fig. 5).

The \( C_{3v} \) site symmetry of the MoO4 ion splits the triply degenerate \( v_{3} \) and \( v_{4} \) vibrations \( (F_{2} \rightarrow A_{1}+E) \). The presence of two molybdate ions in the primitive cell, interacting with each other, results in the additional Davydov (or factor group) splitting, namely, \( v_{3}, v_{4}(F_{2}) \rightarrow (A_{1g}+A_{2g})+(E_{g}+E_{u}) \). Note, that the motions of Rb and Fe atoms do not manifest themselves in the Raman spectra.

**D. Raman spectra**

Figure 6(a) displays unpolarized Raman spectra of RbFe(MoO4)2 at room temperature and at 33 K. Below in this section, the temperatures of the sample holder are indicated and the heating of the sample by laser light is not taken into account. The room-temperature spectrum consists of nine lines, which is consistent with the \( P\bar{3}m1 \) space group.

Table II gives the assignments of Raman modes deduced from the partially polarized spectra [see Fig. 6(b)] and from a comparison with literature data on similar compounds\(^{11} \) and on aqueous solution of \( (\text{MoO}_4)^{2-} \), Ref. 12.

All the Raman lines shift and narrow with decreasing temperature (see Fig. 7). The temperature dependences of the line positions (Fig. 8) and linewidths (Fig. 9) show a specific feature at \( T_{c} \approx 180 \) K that evidences a phase transition. In particular, the distance between the two strong lines \( (A_{1g}+E_{g}) \) originating from \( v_{4}(F_{2}) \) vibration of the \( (\text{MoO}_4)^{2-} \) tetrahedron, grows below the transition temperature, reflecting an additional distortion of the mentioned tetrahedron. The local \( C_{3v} \) symmetry of \( (\text{MoO}_4)^{2-} \) either does not lower or lowers to \( C_{3} \) (no splitting of the doubly degenerate components of the \( v_{4} \) or \( v_{3} \) vibrations is observed).

In the lowest-frequency region of the Raman spectra a new line near \( 80 \) cm\(^{-1} \) appears in the low-temperature spectrum [Figs. 6(a) and 7(a)]. This line strongly broadens and shifts to lower frequencies when approaching \( T_{c} \) from below [see Figs. 7(a), 8(a), and 9(a)].

**IV. DISCUSSION**

The Raman, ESR, and x-ray diffraction data point to a structural phase transition in RbFe(MoO4)2 at \( T_{c} \approx 190 \) K.

**TABLE II.** Raman modes of RbFe(MoO4)2 at room temperature.

<table>
<thead>
<tr>
<th>Frequency (cm(^{-1}))</th>
<th>Symmetry</th>
<th>Assignment</th>
</tr>
</thead>
<tbody>
<tr>
<td>977</td>
<td>( A_{1g} )</td>
<td>( v_{1}(A_{1}) ) vibration of ( \text{MoO}<em>4^{2-} ) ( (T</em>{d}) )</td>
</tr>
<tr>
<td>930</td>
<td>( E_{g} )</td>
<td>( v_{3}(F_{2}) ) vibration of ( \text{MoO}<em>4^{2-} ) ( (T</em>{d}) )</td>
</tr>
<tr>
<td>765</td>
<td>( A_{1g} )</td>
<td>( v_{2}(E) ) vibration of ( \text{MoO}<em>4^{2-} ) ( (T</em>{d}) )</td>
</tr>
<tr>
<td>360</td>
<td>( E_{g} )</td>
<td>( v_{4}(E) ) vibration of ( \text{MoO}<em>4^{2-} ) ( (T</em>{d}) )</td>
</tr>
<tr>
<td>343</td>
<td>( E_{g} )</td>
<td>transnational motions of ( \text{MoO}_4^{2-} )</td>
</tr>
<tr>
<td>168</td>
<td>( A_{1g} )</td>
<td>translational motions of ( \text{MoO}_4^{2-} )</td>
</tr>
<tr>
<td>150</td>
<td>( E_{g} )</td>
<td>librational motion of ( \text{MoO}_4^{2-} )</td>
</tr>
</tbody>
</table>
In view of the temperature $T_c \approx 180$ K that follows from Raman measurements a heating of the sample by laser light should be about 10 K which seems to be reasonable. The Raman spectrum does not change qualitatively below $T_c$, except for the appearance of one new low-frequency mode. Such a mode could be of magnetic origin, due to, e.g., two-magnon Raman scattering. In low-dimensional systems magnetic modes are observed well above $T_N$, due to short-range antiferromagnetic order (see, e.g., Ref. 13). However, the highest frequency in the magnon spectrum of RbFe(MoO$_4$)$_2$ as calculated from the value of the exchange field 67 kOe (Ref. 5) and according to the calculations of the spin-wave spectrum$^{14}$ is only 6.3 cm$^{-1}$, by far too low to account for the 80 cm$^{-1}$ Raman mode. As an alternative interpretation, one could suggest the amplitude mode of an incommensurate structural phase transition. The asymmetric reflections observed in the temperature interval below $T_c$ could be due to the appearance of satellites in an incommensurate phase. In that case, the phase transition at $T_c \approx 190$ K into an incommensurate crystal structure would be a precursor of the incommensurate magnetic phase below $T_N = 3.8$ K. However, in all the investigated incommensurate crystals the frequency of the amplitude mode does not exceed 40 cm$^{-1}$ (see, e.g., Ref. 15), which makes the aforementioned interpretation unlikely.

Double molybdates and tungstates with general formula $A^+R^{3+}(MO_4)_2$ (here $A$ is alkaline atom, Ag, or Tl, $R$ is rare earth, Bi, In, Sc, Ga, Al, Fe, or Cr, $M$ stands for Mo or W) demonstrate a vast variety of crystal structures, exhibit polymorphic transformations and structural phase transitions.$^{1,3,7}$

A thorough crystallographic and x-ray structural study on single crystals of KAl(MoO$_4$)$_2$, KSc(MoO$_4$)$_2$, and KFe(MoO$_4$)$_2$ at room temperature revealed a subtle difference between the structures of the first two and the third compound, in spite of a complete identity of their x-ray powder-diffraction patterns.$^3$ The structures of KR(MoO$_4$)$_2$ with $R$=Al or Sc belong to the $P\overline{3}m1$ space group whereas the structure of KFe(MoO$_4$)$_2$ is described by the $P\overline{3}c1$ space group. The latter structure is characterized by a pseudoperiod that doubles the $c$ parameter. Successive layers are not translationally equivalent, as in the case of the KAl(MoO$_4$)$_2$ structure, but are interconnected by a glide plane $c$. This structural difference is mediated by tiny displacements of the oxygen atoms O2 whereas the positions of other atoms do not change.$^2$ The transformation from $P\overline{3}m1$ to $P\overline{3}c1$ may be considered as a rotation of the two (MoO$_4$)$_2^-$ tetrahedra in the primitive cell of the $P\overline{3}m1$ structure in opposite directions about the $C_3$ axis. As the O2 oxygen atoms belong also to the FeO$_6$ octahedra, these rotations lead to the appearance of two differently oriented octahedra in the new $P\overline{3}c1$ primitive cell.

It is possible that the structural phase transition in RbFe(MoO$_4$)$_2$ at $T_c \approx 190$ K is just a transformation from a more symmetric $P\overline{3}m1$ room-temperature structure into a less symmetric $P\overline{3}c1$ low-temperature one, that comprises two differently oriented FeO$_6$ octahedra.

Table III gives atomic positions and the appropriate irre-
tributions to the vibrational modes of RbFe(MoO$_4$)$_2$ in the case of $E_u$ vibrations of the MoO$_4$ tetrahedron, due to the additional Davy-

There are 16 Raman $(5A_{1g}+11E_g)$ and 18 infrared $(6A_{2u}+12E_u)$ modes. Analyzing the vibrational modes in more detail we find that four additional $E_u$ modes and a new doublet $E_g+A_{1g}$ should appear in the region of $\nu_2$, $\nu_3$, $\nu_4$ vibrational, translational and, respectively, librational motions of the MoO$_4$ tetrahedron, due to the additional Davy-

\begin{equation}
\Gamma^{opt}(P\bar{3}c1)=5A_{1g}(xx,yy,zz)+11E_g(xz,yz,xy)+6A_{2g}+6A_{1u}+6A_{2u}(E||z)+12E_u(E||x,y).\nonumber
\end{equation}

TABLE III. Crystallographic positions and corresponding contributions to the vibrational modes of RbFe(MoO$_4$)$_2$ in the case of $P\bar{3}c1$ structure.

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<td>$D_3$</td>
<td>$A_{2g}+A_{2u}+E_g+E_u$</td>
</tr>
<tr>
<td>Fe</td>
<td>2b</td>
<td>$C_{3i}$</td>
<td>$A_{1u}+A_{2u}+2E_u$</td>
</tr>
<tr>
<td>Mo</td>
<td>4d</td>
<td>$C_3$</td>
<td>$A_{1g}+A_{1u}+A_{2g}+A_{2u}+2E_g+2E_u$</td>
</tr>
<tr>
<td>O1</td>
<td>4d</td>
<td>$C_3$</td>
<td>$A_{1g}+A_{1u}+A_{2g}+A_{2u}+2E_g+2E_u$</td>
</tr>
<tr>
<td>O2</td>
<td>12g</td>
<td>$C_1$</td>
<td>$3A_{1g}+3A_{1u}+3A_{2g}+3A_{2u}+6E_g+6E_u$</td>
</tr>
</tbody>
</table>

dov splitting in the doubled unit cell. The sixth new $E_g$ mode comes from an activation of Rb motions in the Raman spectra.

The additional Davydov splitting can be small because of a weak interaction between MoO$_4$ tetrahedra in different layers, so that it is not resolved in Raman spectra. The exception is the additional $A_{1g}$ mode, a Davydov counterpart to the silent $A_{2g}$ librational mode of the $P\bar{3}m1$ structure. We suggest that the extra Raman mode observed below $T_c$ in RbFe(MoO$_4$)$_2$ is just this new $A_{1g}$ librational mode. We have to admit, however, that a new $E_u$ mode associated with motions of Rb is too weak to be observed.

The ESR data enable us to estimate the change of the crystal field below the phase-transition point $T_c=190$ K, by means of the following considerations. The temperature dependence of the ESR intensity above 10 K corresponds well to the Curie-Weiss law $(T+\Theta_{CW})^{-1}$ of the static susceptibility with the Curie-Weiss constant $\Theta_{CW}\approx25$ K. This value gives an estimation of the exchange interaction energy. Thus, in the temperature range above 100 K, far from critical broadening due to the onset of magnetic order, we consider the ESR linewidth in the high-temperature approximation $(T>>\Theta_{CW})$. To show that the crystal field is responsible for the line broadening in RbFe(MoO$_4$)$_2$, we first determine its strength from the evaluation of the resonance field measured at 34 GHz.

The temperature dependence of the resonance field at 34 GHz exhibits a signature typical for the influence of an internal anisotropy field, where the resonance shift is related to the Curie-Weiss law of the susceptibility. The symmetry of the crystal field at the iron site allows only for a uniaxial spin anisotropy with the anisotropy axis parallel to the crystallographic $c$ axis. We choose this axis as the $z$ axis of the local coordinate system. The crystal field partially lifts the degeneracy of the Zeeman levels of the spin-$S=5/2$ ion Fe$^{3+}$, already without magnetic field, and leads to a splitting of the ESR spectrum. This effect is usually described in terms of a spin Hamiltonian with a single-ion anisotropy term $D_{z^2}$, where $D$ is the anisotropy constant.\(^{16}\) However, the splitting is not observed in RbFe(MoO$_4$)$_2$, because the influence of the crystal field is effectively averaged by the isotropic Heisenberg exchange interaction. The ESR spectrum is exchange narrowed into a single resonance line, with line shift and broadening resulting from the underlying unnarrowed spectrum.

From the equation of motion we derive the following expressions describing the influence of the crystal field on the temperature dependence of the resonance field $H_{res}$, analogously to Ref. 17:

\begin{equation}
H_{res}(c\parallel H)=\frac{h\nu}{g\mu_B}\frac{1-2S(S+1)D}{3k_B(T-\Theta_{CW})}^{-1},
\end{equation}

\begin{equation}
H_{res}(c\perp H)=\frac{h\nu}{g\mu_B}\frac{1+2S(S+1)D}{3k_B(T-\Theta_{CW})}^{-1/2}.
\end{equation}

Here $\nu$ is the microwave frequency, $h$ is the Planck constant, $\mu_B$ is the Bohr magneton, and $k_B$ is the Boltzmann
cally ordered phase 5, where the anisotropy field,
et and ESR measurements in the antiferromagnetic
very well with those obtained independently from the mag-
the form


p

The corresponding fit curves are shown as
solid lines in the lower frame of Fig. 4. The value
is nearly isotropic, as expected for S-state ions. The Curie-Weiss
temperatur is in good agreement with that found from the mag-
netic susceptibility measurements.

Having determined the anisotropy constant D, we can
roughly estimate the expected crystal-field contribution to
the line width, in the framework of the theory of exchange
narrowing, 18,19 by

\[
\Delta H_{\text{CF}} = \frac{1}{g \mu_B} \frac{D^2}{J}.
\]

(5)

With the Heisenberg exchange coupling \( J \), of a spin \( S \),
to its six nearest neighbors \( S \), in the basal plane, we obtain
the corresponding exchange integral from the Curie-Weiss
temperature as \( J/k_B = \Theta_{\text{CW}}/(6S^2) = 0.67 \) K. Inserting
the values for \( D \) and \( J \) in Eq. (5) yields \( \Delta H_{\text{CF}} = 700 \) Oe, which
is the right order of magnitude.

Note that the values of \( D \) and \( J \) determined above agree
very well with those obtained independently from the mag-
netization and ESR measurements in the antiferromagneti-
cally ordered phase 5, where the anisotropy field \( H_0 = DS/(g \mu_B) = 5 \) kOe and the exchange field \( H_E = 6JS/(g \mu_B) = 67 \) kOe give \( D/k_B = 0.27 \) K and \( J/k_B = 0.6 \) K, respectively.

Concerning other possible contributions to the linewidth,
symmetric anisotropic exchange and Dzyaloshinsky-Moriya
interaction can be neglected. Indeed, the orders of magni-
tudes of the characteristic energies of these interactions are
(\( \Delta g/g \))\( J \) and (\( \Delta g/g \))\( J \), respectively, where (\( \Delta g/g \)) is the
relative anisotropy of the g factor. 20 Due to the nearly isotro-
pic g factor observed in our measurements the magnitudes of
the interactions mentioned above are negligible in comparison
with \( D \). Hyperfine interaction can be neglected as well,
because the nuclear spins of \( ^{56}\text{Fe} \) and surrounding \( ^{16}\text{O} \)
are zero and the other nuclei are too far away. Only the dipole-
dipole interaction gives an additional contribution in the
presence of the exchange interaction, and can be written in the
form

\[
\Delta H_{\text{DD}} = \frac{3g^3 \mu_B^3 S(S+1)}{J r^6}.
\]

(6)

With the nearest-neighbor distance \( r = a = 5.69 \) \( \text{Å} \), this
yields \( \Delta H_{\text{DD}} = 50 \) Oe, which is still small compared to the experimentally observed linewidth. Hence the linewidth is
indeed dominated by the single-ion anisotropy. This is also in
agreement with the observed angular dependence of the line-
width. Namely, for the magnetic field applied parallel to the
direction of the anisotropy axis the ESR line is twice as large
as for the magnetic field applied perpendicular to this direc-
tion.

Transformation of the crystal structure causes in a natural
way a change of the crystal field in the surrounding of Fe\(^{3+}\)
ions. Using the above consideration of the ESR linewidth we
deduce from the 4% change of the linewidth in the tempera-
ture range between 180 and 190 K that the crystal field is
increased by about 2% below the transition point. The in-
crease of the crystal field strength may be caused by the
distortion of the FeO\(_6\) octahedra combined with the rotations
of MoO\(_4\), because they are sharing vertices.

V. CONCLUSIONS

We have carried out x-ray powder diffraction, ESR, and
Raman scattering measurements on RbFe(MoO\(_4\))\(_2\) at different
temperatures. The temperature dependence of the unit-
cell parameters \( a \) and \( c \) display a peculiarity between 200 and
180 K. The temperature dependences of the resonance fields
demonstrate small but distinct steps at \( T_c \approx 190 \) K. In the
Raman spectra, all nine modes predicted by the group-
theoretical analysis were observed for the temperatures
higher than \( T_c \). Below \( T_c \), these modes shift and narrow
markedly and a new low-frequency mode appears. Raman,
ESR, and x-ray powder-diffraction data point to a second
order structural phase transition at \( T_c \approx 190 \) K. Based on the changes in the Raman and ESR spectra below \( T_c \), we sug-
gest that the 190 K phase transition is mediated by rotations
of the MoO\(_4\) tetrahedra, resulting in a symmetry lowering
from \( P\bar{3}m1 \) to \( P\bar{3}c1 \). From the experimental temperature
dependencies of the ESR field above 10 K we derived the
\( g \)-factor values \( g_\| = 1.992 \) and \( g_\perp = 2.003 \), the Curie-Weiss
constant \( \Theta_{\text{CW}} = -21 \) K, and the single-ion anisotropy con-
stant \( D/k_B = 0.25 \) K.

ACKNOWLEDGMENT

The authors kindly acknowledge D. Vieweg for making
x-ray and SQUID measurements, and V.N. Glazkov and S.S.
Sosin for discussions. This work was supported by INTAS,
Grant No. 99–0155 and by the Russian Foundation for Basic
Research, Grants No. 01–02–16329 and 01-02-17557, by
Deutsche Forschungsgemeinschaft (DFG) via Grant No. SFB
484 and the joint Grant Nos. DFG-RFBR, 436 RUS 113/628,
by German BMBF under Contract No. VDI/EMK 13N6917,
and by the Netherlands Foundation for Fundamental Re-
search on Matter (FOM).

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