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LETTER TO THE EDITOR

## Detection of HF emission from the Orion Bar <sup>★</sup>

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### ABSTRACT

**Context.** The clumpy density structure of photon-dominated regions is well established, but the physical properties of the clumps and of the surrounding interclump medium are only approximately known.

**Aims.** The aim of this paper is to constrain the physical and chemical conditions in the Orion Bar, a prototypical nearby photon-dominated region.

**Methods.** We present observations of the HF  $J=1-0$  line, which appears in emission toward the Orion Bar, and compare the brightness of the line to non-LTE radiative transfer calculations.

**Results.** The large width of the HF line suggests an origin of the emission in the interclump gas, but collisional excitation by  $H_2$  in the interclump gas underpredicts the observed line intensity by factors of 3–5. In contrast, an origin of the line in the dense clumps requires a density of  $\sim 10^9 \text{ cm}^{-3}$ , 10–100 times higher than previous estimates, which is unlikely. However, electron impact excitation reproduces our observations for  $T=100 \text{ K}$  and  $n_e=10 \text{ cm}^{-3}$ , as expected for the interclump gas.

**Conclusions.** We conclude that HF emission is a signpost of molecular gas with a high electron density. Similar conditions may apply to active galactic nuclei where HF also appears in emission.

**Key words.** ISM: molecules – astrochemistry

## 1. Introduction

Photon-dominated regions (PDRs) are the surface regions of molecular clouds, where ultraviolet radiation with photon energies between a few and 13.6 eV drives the thermal and chemical balance of the gas (Hollenbach & Tielens 1999). This situation occurs in regions of high-mass star formation, but also in protoplanetary disks and in the nuclei of active galaxies. Studying the structure of PDRs therefore has a wide astrophysical application. In PDRs, gas heating proceeds by photo-electric emission from dust grains, while the main cooling channels are the fine structure lines of  $C^+$  and O and the rotational lines of CO (Kaufman et al. 1999). Absorption of the impinging ultraviolet radiation by dust and gas in the PDR creates a layered structure, where chemical transitions such as  $H^+ \rightarrow H \rightarrow H_2$  and  $C^+ \rightarrow C \rightarrow CO$  occur.

The Orion Bar is a prototypical PDR, located between the Orion molecular cloud and the Orion Nebula, the  $H^+$  region surrounding the Trapezium stars, at a distance of 414 pc (Menten et al. 2007). Observations at infrared and submillimeter wavelengths indicate a geometry for the Bar where the PDR is wrapped around the Orion Nebula, and changes from a face-on to an edge-on view where the molecular emission peaks (Hogerheijde et al. 1995; Walmsley et al. 2000). The mean temperature of the molecular gas in the Bar is 85 K, while the temperature rises to several 100 K toward the ionization front, where

the emission from PAH particles and vibrationally excited  $H_2$  peaks.

While the temperature structure of the Orion Bar is reasonably well understood, the same cannot be said about the density structure. The mean density of the molecular gas is  $10^5 \text{ cm}^{-3}$ , but single-dish observations already indicate the presence of small-scale density variations without apparent pattern, usually called "clumps" (Hogerheijde et al. 1995), that are also seen toward other PDRs (Stutzki et al. 1988; Wang et al. 1993). While interferometric observations have confirmed the presence of clumps (Young Owl et al. 2000), the densities of both the clumps and the interclump medium are somewhat uncertain. The interclump medium probably has a density between a few  $10^4$  and  $2 \times 10^5 \text{ cm}^{-3}$  (Simon et al. 1997), while estimates of the clump density range from  $1.5 \times 10^6$  to  $6 \times 10^6 \text{ cm}^{-3}$  (Lis & Schilke 2003).

This Letter presents observations of the HF  $J=1-0$  line, which appears in emission toward the Orion Bar. The HF molecule is expected to be the dominant carrier of gas-phase fluorine, because the reaction  $F + H_2 \rightarrow HF + H$  is exothermic. For diffuse clouds where the effect of depletion on grains should be unimportant, models by Neufeld & Wolfire (2009) predict an HF abundance of  $\sim 3.6 \times 10^{-8}$  relative to  $H_2$ . Recent observations of the HF  $1-0$  line confirm this prediction: the line is seen in absorption toward several background sources, indicating abundances of  $\sim 1-2 \times 10^{-8}$  (Neufeld et al. 2010). Toward dense clouds, the abundance is measured to be  $\sim 100$  times lower (Phillips et al. 2010), suggesting that depletion of F on grain surfaces or excitation effects play a role.

<sup>★</sup> *Herschel* is an ESA space observatory with science instruments provided by European-led Principal Investigator consortia and with important participation from NASA

**Table 1.** Line parameters.

$\int T_{\text{mb}} \Delta V$ K km s <sup>-1</sup>	$V_{\text{LSR}}$ km s <sup>-1</sup>	$\Delta V$ km s <sup>-1</sup>	$T_{\text{mb}}$ K
7.79(21)	9.23(7)	4.86(15)	1.50

**Notes.** Numbers in parentheses are error bars in units of the last decimal.

Like CO, HF is a linear rotor with a regular line spectrum, where  $\nu = 2B(J + 1)$  and  $A_{ij} \propto \nu^3$ . Unlike CO, HF has a small reduced mass and a large dipole moment, so that the lines have high frequencies and radiative decay is rapid. In particular, the HF 1–0 line has a frequency of 1232 GHz and an Einstein A coefficient of  $2.422 \times 10^{-2} \text{ s}^{-1}$ . Thermal excitation of the  $J=1$  level of HF thus requires extremely high gas densities, which is why the line usually appears in absorption (Sonnentrucker et al. 2010; Monje et al. 2011a).

## 2. Observations and data reduction

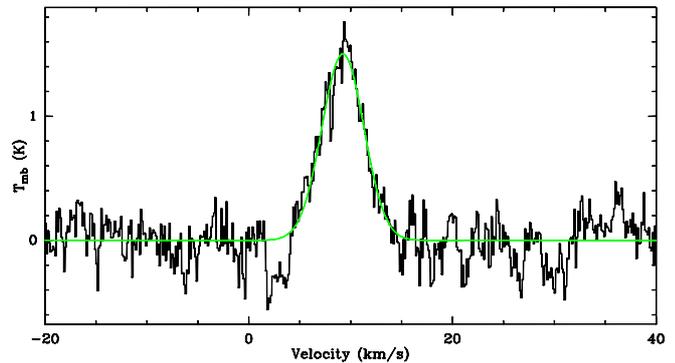
The observations were made with the Heterodyne Instrument for the Far-Infrared (HIFI; De Graauw et al. 2010) onboard ESA’s *Herschel* Space Observatory (Pilbratt et al. 2010) on 2011 March 19, as part of the Guaranteed Time key program *Herschel* Observations of Extraordinary Sources (HEXOS; Bergin et al. 2010). Receiver band 5b was used as front end and the acousto-optical Wide-Band Spectrometer (WBS), which covers 4 GHz bandwidth in four 1140 MHz subbands at 1.1 MHz (0.30 km s<sup>-1</sup>) resolution, as back end. The data were taken in spectral survey mode with a redundancy of 6, using the frequency switch technique with a throw of 100 MHz.

The position observed is the CO<sup>+</sup> peak of the Orion Bar at  $\alpha = 05^{\text{h}} 35^{\text{m}} 20.6^{\text{s}}$ ,  $\delta = -05^{\circ} 25' 14''$  (J2000) (Störzer et al. 1995). The FWHM beam size at the observing frequency is 17.3 arcsec (Roelfsema et al. 2011), which corresponds to 7200 AU at the distance of Orion. The system temperature for these data is 2571 K in SSB main beam units, and the integration time is 70 minutes (ON+OFF).

Calibration of the data, removal of standing waves, and side-band deconvolution were performed in the *Herschel* Interactive Processing Environment (HIPE) version 6.0. The subsequent analysis was performed within the CLASS<sup>1</sup> package, version January 2011. The calibration is estimated to be accurate to  $\approx 10\%$ ; the velocity scale should be accurate to 0.5 km s<sup>-1</sup> or better. The intensity scale was converted to  $T_{\text{mb}}$  using a main beam efficiency of 64% (Roelfsema et al. 2011). Only a linear baseline was subtracted. After inspection, data from the two polarizations were averaged together to obtain an rms noise level of 0.18 K in  $T_{\text{mb}}$  units on 0.5 MHz channels.

## 3. Results

Figure 1 presents the HF line profile observed toward the Orion Bar. Strong emission is detected at  $V_{\text{LSR}} = +9.24 \text{ km s}^{-1}$ , with an almost Gaussian shape. In addition, weak absorption may be present near  $V_{\text{LSR}} = +2..+4 \text{ km s}^{-1}$ , which appears in both polarizations. If confirmed, this absorption arises in the Orion Ridge (Tauber et al. 1994). Standing waves prohibit the detection of continuum signal in our data down to an rms of  $\approx 1 \text{ K}$  in  $T_{\text{mb}}^*$ .



**Fig. 1.** Spectrum of the HF  $J=1-0$  line, observed with *Herschel*-HIFI toward the Orion Bar.

Table 1 lists the parameters of the emission feature, measured by fitting a Gaussian profile to the data in Figure 1. The measured  $V_{\text{LSR}}$  agrees well with values for the dense molecular gas in the Orion Bar as probed by H<sub>2</sub>CO and CH<sub>3</sub>OH lines in the 100–400 GHz range (Hogerheijde et al. 1995; Leurini et al. 2006). However, the FWHM line width of 4.86 km s<sup>-1</sup> is much broader than measured for the dense molecular gas, by a factor of  $\geq 2$ . This difference is about twice as large as expected from the lower mass of the HF molecule compared with species such as CO. The width of the HF line is close to the values measured for C<sup>+</sup> and CH<sup>+</sup> (Nagy et al, in prep.), suggesting an origin of the observed HF emission in the low-density interclump gas of the Orion Bar, which is a surprise given the large critical density ( $\sim 5 \times 10^9 \text{ cm}^{-3}$ ) of the line.

## 4. Excitation by H<sub>2</sub>

To clarify the origin of the observed HF emission, we have performed model calculations with the RADEX program (Van der Tak et al. 2007). This program solves the equations of statistical equilibrium in the presence of collisional and radiative excitation, treating optical depth effects in the lines with an escape probability formalism. The calculations use spectroscopic input data from the JPL catalog (Pickett et al. 1998), and collisional rate coefficients for the HF-H<sub>2</sub> system from Guillon & Stoecklin (2011). These collision rates are  $\approx 10$  times larger than calculated for the HF-He system (Reese et al. 2005), even after scaling for the difference in reduced mass (Schöier et al. 2005). This difference is at the upper end of the expected range (Van der Tak 2011), due to the large well depth of the HF-H<sub>2</sub> interaction potential. All calculations assume a line width of 5 km s<sup>-1</sup> (Table 1) and a background radiation field of 2.73 K due to the cosmic microwave background. We assume that the ortho-para ratio of H<sub>2</sub> equals its thermal value. The free parameters in the calculation are the kinetic temperature, the gas density  $n(\text{H}_2)$  and the HF column density  $N(\text{HF})$ .

Based on the measured line width, we first consider an origin of the HF emission in the warm interclump gas close to the ionization front. We adopt  $n(\text{H}_2) = 10^5 \text{ cm}^{-3}$ ,  $N(\text{H}_2) = 10^{22} \text{ cm}^{-2}$  (Hogerheijde et al. 1995), and assume that all fluorine (at Solar abundance) is in the form of HF, as is commonly observed for

<sup>1</sup> <http://www.iram.fr/IRAMFR/GILDAS>

diffuse interstellar clouds. This assumption implies an HF abundance of  $3.6 \times 10^{-8}$ , so that  $N(\text{HF}) = 3.6 \times 10^{14} \text{ cm}^{-2}$ . For a gas kinetic temperature of 100 K, this model predicts an intensity for the HF 1–0 line of 0.5 K, which is a factor of 3 lower than observed. Models in which the HF emission arises in warm surface layers with a low molecular fraction and accordingly a low HF abundance also do not match our observations, provided that HF-H collision rates are similar to those for HF-H<sub>2</sub>. We conclude that the observed HF emission does not originate in the warm interclump gas if collisional excitation by H<sub>2</sub> or H is responsible.

As an alternative model, we consider an origin of the HF emission in the dense clumps of the Orion Bar. For this gas, we assume a kinetic temperature of 85 K (Hogerheijde et al. 1995),  $N(\text{H}_2) \approx 1 \times 10^{23}$  (e.g., Van der Wiel et al. 2009) and an HF abundance of  $\lesssim 3.6 \times 10^{-8}$ . These models reproduce the observed HF line intensity for gas densities in the range  $10^8$ – $10^9 \text{ cm}^{-3}$  and  $N(\text{HF})$  in the range  $10^{12}$ – $10^{13} \text{ cm}^{-2}$ . The line optical depth in these models is around unity, which is consistent with the observed line shape. The implied HF abundance is  $\sim 10^{-10}$ – $10^{-11}$ , which is similar to the values found for dense gas in previous studies (Phillips et al. 2010). Since the line is effectively optically thin, lower H<sub>2</sub> volume densities require correspondingly higher HF column densities, so that the elemental fluorine abundance would be exceeded. We conclude that thermal excitation of the HF line indicates a high gas density of  $n = 10^8$ – $10^9 \text{ cm}^{-3}$ , similar to the inner envelopes of late-type stars where HF emission is also seen (Agúndez et al. 2011).

One problem with this alternative model is that the required gas density of  $\sim 10^9 \text{ cm}^{-3}$  is much higher than commonly adopted for the Orion Bar. An origin of the HF emission in high-density clumps is also inconsistent with the line width, which is much larger than that of CO. In addition, while peak densities of  $\sim 10^9 \text{ cm}^{-3}$  may occur locally, this value is unreasonably high as an average over the 7200 AU HIFI beam. Finally, the Orion Bar does not stand out for its high density between regions where HF has been observed, which all show the line in absorption.

## 5. Excitation by electrons

Although collisional excitation of HF by H<sub>2</sub> does not match our observations of the Orion Bar, H<sub>2</sub> is not the only possible collision partner responsible for the excitation of HF in PDR-type regions. In the interclump gas of the Orion Bar, most carbon is in the form of C<sup>+</sup>, leading to an electron abundance of  $\sim 10^{-4}$  and an electron density of  $10 \text{ cm}^{-3}$ . Under these conditions, a model without collisional excitation by electrons would be incomplete. Quantum mechanical calculations of collisional cross sections for the e-HF system only exist for  $T > 500 \text{ K}$  (Thümmel et al. 1992), but for highly polar molecules such as HF, the Born approximation (see, e.g., Itikawa & Mason 2005) should give reliable results as long as  $kT > E_{\text{up}}$ . We have calculated electron impact excitation rates for the  $\Delta J = 1$  transitions of HF in the Born approximation, and scaled the values to the results of the R-matrix calculations by Thümmel et al. (1992). These rates should be accurate to within a factor of  $\approx 2$ , similar to the well-studied e-H<sub>2</sub>O system with a similar dipole moment (Zhang et al. 2009), and sufficient for our purposes. Calculations with RADEX indicate that for  $T = 100 \text{ K}$  and  $n_e = 10 \text{ cm}^{-3}$ , the observed HF line intensity is reproduced for  $N(\text{HF}) \approx 1 \times 10^{15} \text{ cm}^{-2}$ .

The corresponding H<sub>2</sub> column density may be estimated from the C<sup>+</sup> observations by Herrmann et al. (1997) with the KAO in a 55'' beam. These data indicate that  $N(\text{C}^+) \approx (3\text{--}5) \times 10^{18} \text{ cm}^{-2}$ , so that for  $\text{C}^+/\text{H}_2 = 2 \times \text{C}/\text{H} = 2.4 \times 10^{-4}$

(Wakelam & Herbst 2008; Sofia et al. 2011), we obtain  $N(\text{H}_2) = (1.3\text{--}2.0) \times 10^{22} \text{ cm}^{-2}$ . This value is a lower limit because of beam dilution: mapping of C<sup>+</sup> with HIFI indicates that  $N(\text{C}^+)$  and therefore  $N(\text{H}_2)$  is  $\sim 3\text{--}5$  times higher (Röllig et al. in prep.). The HF abundance in the C<sup>+</sup>-H<sub>2</sub> layer is thus  $(3\text{--}4) \times 10^{-8}$ , consistent with the Solar fluorine abundance. We conclude that electron excitation in the interclump medium is a viable model to explain the observed HF emission in the Orion Bar.

## 6. Conclusions

We have presented observations of the HF  $J=1\text{--}0$  line toward the Orion Bar, which is the first time that this line is seen in emission from the Galactic interstellar medium. We present calculations and arguments why this emission cannot be caused by collisional excitation of HF by H<sub>2</sub>, either from dense clumps or from the inter-clump medium. However, collisional excitation by electrons in the interclump gas explains our observations. The appearance of HF in emission therefore seems to be a signpost of molecular gas with a high electron density, which in the case of the Orion Bar is the combined effect of the high gas density and the strong ultraviolet radiation field.

The physical conditions of the Orion Bar may be similar to those in the nucleus of the active galaxy Mrk 231, where the radiation field is strong and where the HF line also appears in emission (Van der Werf et al. 2010). Other active galactic nuclei where HF appears in absorption such as Arp 220 (Rangwala et al. 2011) and the Cloverleaf (Monje et al. 2011b) probably have lower electron densities, which may indicate a softer radiation field.

Support for an origin of the HF emission in the interclump medium comes from the detection of CF<sup>+</sup> toward the Orion Bar (Neufeld et al. 2006). Additional support for this model would come from observations of H<sub>2</sub>F<sup>+</sup>, for which line frequencies are known (Fujimori et al. 2011). However, our data do not show any features at the predicted line frequencies down to an rms of  $\approx 50 \text{ mK}$  in  $T_A^*$ , which is a factor of  $\sim 10$  above the expected signal.

In the future, maps of HF emission with Herschel-HIFI will help to clarify the spatial distribution of dense gas clumps in the Orion Bar and other PDRs. Another useful test would be the observation of HF  $J=2\text{--}1$  or higher- $J$  lines in the  $v=0$  or  $v=1$  states. Unfortunately, the  $J=2\text{--}1$  line at 2463 GHz lies very close ( $0.13 \mu\text{m}$ ) to the N<sup>+</sup>  $^3P_2\text{--}^3P_1$  fine-structure line at  $121.8 \mu\text{m}$ , so that the lines are blended at the resolution of the PACS instrument onboard Herschel. Heterodyne observations in this frequency range may become possible with the future STO telescope, the SOFIA airborne observatory, or the Millimetron mission.

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