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First Study of Antihyperon-Nucleon Scattering $\bar{\Lambda}p \to \bar{\Lambda}p$ and Measurement of $\Lambda p \to \Lambda p$ Cross Section

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Using $(10.087 \pm 0.044) \times 10^9$ J/ψ events collected with the BESIII detector at the BEPCII storage ring, the processes $\Lambda p \to \Lambda p$ and $\bar{\Lambda}p \to \bar{\Lambda}p$ are studied, where the $\Lambda/\bar{\Lambda}$ baryons are produced in the process $J/\psi \to \Lambda\bar{\Lambda}$ and the protons are the hydrogen nuclei in the cooling oil of the beam pipe. Clear signals are observed for the two reactions. The cross sections in $-0.9 \leq \cos\theta_{\Lambda/\bar{\Lambda}} \leq 0.9$ are measured to be $\sigma(\Lambda p \to \Lambda p) = (12.2 \pm 1.6_{\mathrm{stat}} \pm 1.1_{\mathrm{syst}})$ and $\sigma(\bar{\Lambda}p \to \bar{\Lambda}p) = (17.5 \pm 2.1_{\mathrm{stat}} \pm 1.6_{\mathrm{syst}})$ mb at the $\Lambda/\bar{\Lambda}$ momentum of 1.074 GeV$/c$ within a range of ± 0.017 GeV$/c$, where the $\theta_{\Lambda/\bar{\Lambda}}$ are the scattering angles of the $\Lambda/\bar{\Lambda}$ in the $\Lambda p/\bar{\Lambda}p$ rest frames. Furthermore, the differential cross sections of the two reactions are also measured, where there is a slight tendency of forward scattering for $\Lambda p \to \Lambda p$, and a strong forward peak for $\bar{\Lambda}p \to \bar{\Lambda}p$. We present an approach to extract the total elastic cross sections by extrapolation. The study of $\bar{\Lambda}p \to \bar{\Lambda}p$ represents the first study of antihyperon-nucleon scattering, and these new measurements will serve as important inputs for the theoretical understanding of the (anti)hyperon-nucleon interaction.

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One of the main goals of nuclear physics is to understand baryon-baryon interaction in a unified perspective. To achieve this purpose, plentiful nucleon-nucleon (NN) and antinucleon-nucleon ($\bar{\mathrm{N}}$N) scattering data have been measured [1]. Therefore, the relevant theory of NN and $\bar{\mathrm{N}}$N interactions is well established, and it can be tightly constrained by experimental data. However, the understanding of hyperon-nucleon (YN) interaction has a large uncertainty due to the lack of relevant measurements. The YN interaction is studied mainly via three methods. The first is to extract the YN correlation functions in heavy-ion collisions [2–5], the second is to study hypernuclei [6–9], and the third is to investigate YN scattering [10–12]. The last method is the most direct way to study YN interaction, but it is limited by the availability and short-lifetime of hyperon beams, leading to a scarcity of YN scattering data [1]. The study of YN interaction is also crucial to determine the equation of state (EOS) of nuclear matter at supersaturation densities and understand the so-called "hyperon puzzle" of neutron stars (NS) [13–18]. To solve these issues, more YN scattering data are desired to constrain the calculations of YN interaction.

Compared to the YN scattering, the situation is even worse for antihyperon-nucleon ($\bar{\mathrm{Y}}$N) scattering. Until now, no $\bar{\mathrm{Y}}$N scattering data have been obtained due to the absence effective antihyperon sources [1], which results in the very limited related theoretical research. Therefore, the realization of $\bar{\mathrm{Y}}$N scattering measurements can fill this gap, and new measurements will motivate more effort for the understanding of the $\bar{\mathrm{Y}}$N interaction. More importantly, $\bar{\mathrm{Y}}$N scattering data can further constrain the YN interaction theory from another angle.

In this Letter, we present a study of the reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda}p \to \bar{\Lambda}p$, where Λ and $\bar{\Lambda}$ are reconstructed via the decays $\Lambda \to p\pi^-$ and $\bar{\Lambda} \to \bar{p}\pi^+$. The cross sections and differential cross sections of the two reactions are all measured. This is the first study of $\bar{\mathrm{Y}}$N scattering.

The BESIII detector records symmetric e^+e^- collisions at the BEPCII collider [19]. Details of the BESIII detector can be found in Ref. [20]. The material of the beam pipe is composed of gold (^{197}Au), beryllium (^{9}Be), and oil ($^{12}\mathrm{C}:{}^{1}\mathrm{H} = 1:2.13$), as shown in Fig. 1. With a sample of $(10.087 \pm 0.044) \times 10^9$ J/ψ events collected by the BESIII detector [21], intense almost monoenergetic $\Lambda/\bar{\Lambda}$ hyperons with a momentum of 1.074 GeV$/c$ within a range of ± 0.017 GeV$/c$ can be produced via the decay $J/\psi \to \Lambda\bar{\Lambda}$, the momentum spread is due to the small horizontal crossing angle of ± 11 mrad for $e^{\pm}$ beams. Afterwards the $\Lambda/\bar{\Lambda}$ baryons can interact with the material in the beam pipe. A similar idea was proposed forty years ago using $\bar{p}p$ collisions at a LEAR experiment [22]. Especially, Ref. [23] has used this method to perform the

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[*] Full author list given at the end of the Letter.

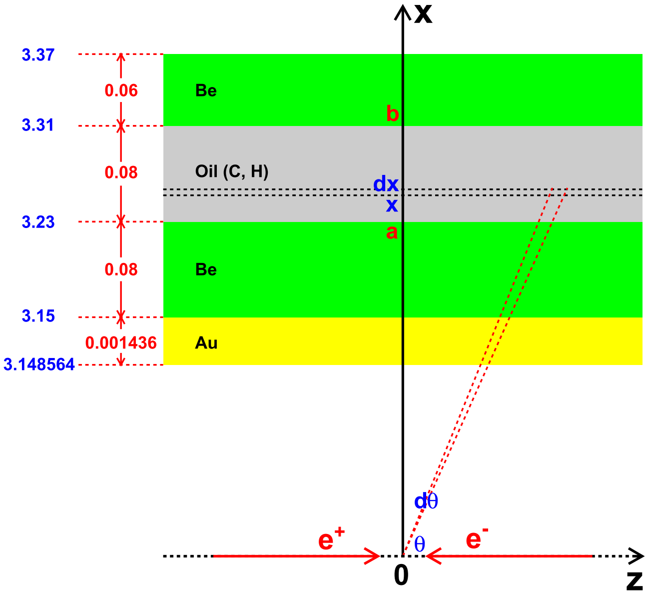

FIG. 1. Schematic diagram of the beam pipe, the length units are centimeter (cm). The z axis is the symmetry axis of the MDC, and the x axis is perpendicular to the e^+e^- beam direction.

first study of YN interaction using Ξ^0-nucleus scattering at BESIII, and Λ-nucleus scattering was measured in Ref. [24]. Furthermore, utilizing the almost static protons in the ^{1}H of the cooling oil of the beam pipe, the information on the interaction between (anti)hyperon and proton can be directly extracted via (anti)hyperon-proton scattering in this way.

In this analysis, simulated data samples are produced with a GEANT4-based [25] Monte Carlo (MC) package, which includes the geometric description of the BESIII detector [20] and the detector response. They are used to determine detection efficiencies and to estimate backgrounds. The simulation models the beam energy spread and initial state radiation (ISR) in the e^+e^- annihilations with the generator KKMC [26]. The inclusive MC sample includes both the production of the J/ψ resonance and the continuum processes incorporated in KKMC [26]. All particle decays are modeled with EVTGEN [27] using branching fractions either taken from the Particle Data Group (PDG) [1], where available, or otherwise estimated with LUNDCHARM [28]. Final state radiation (FSR) from charged final state particles is incorporated using the PHOTOS package [29]. The signal process considered in this analysis is $J/\psi \to \Lambda\bar{\Lambda}$ with either $\Lambda p \to \Lambda p$ or $\bar{\Lambda} p \to \bar{\Lambda} p$, $\Lambda \to p\pi^-$, $\bar{\Lambda} \to \bar{p}\pi^+$. In the signal simulation, the angular distribution of $J/\psi \to \Lambda\bar{\Lambda}$ is generated according to the measurement in Ref. [30]. We simulate the reactions $\Lambda p \to \Lambda p/\bar{\Lambda} p \to \bar{\Lambda} p$ by taking the proton to be at rest, and the hyperon angular distribution is generated using an isotropic phase-space distribution to obtain the angle dependent detection efficiency.

Charged tracks detected in the multilayer drift chamber (MDC) are required to be within a polar angle (θ) range of $|\cos\theta| < 0.93$, where θ is the angle between the charged track and the z axis, which is the symmetry axis of the MDC. Particle identification for charged tracks combines measurements of the energy loss (dE/dx) in the MDC and the flight time in the time-of-flight system (TOF) to form likelihoods $\mathcal{L}(h)$ ($h = p, K, \pi$) for each hadron h hypothesis. Tracks are identified as protons when the proton hypothesis has the greatest likelihood [$\mathcal{L}(p) > \mathcal{L}(\pi)$ and $\mathcal{L}(p) > \mathcal{L}(K)$], while charged pions are identified by comparing the likelihoods for the pion hypotheses, [$\mathcal{L}(\pi) > \mathcal{L}(K)$ and $\mathcal{L}(\pi) > \mathcal{L}(p)$].

Since the final states of the two reactions all contain $pp\bar{p}\pi^+\pi^-$, candidate events must have five charged tracks, and two p, one $\bar{p}$, one π^+, and one π^- are required to be identified. For the decay $\bar{\Lambda} \to \bar{p}\pi^+$, we perform a vertex fit to the $\bar{p}\pi^+$ combination, and the $\bar{\Lambda}$ signal region is defined as $|M(\bar{p}\pi^+) - m_{\bar{\Lambda}}| < 0.003$ GeV/c^2, where $m_{\bar{\Lambda}}$ is the nominal mass of the $\bar{\Lambda}$. In this Letter, all nominal masses are taken from PDG [1]. For the decay $\Lambda \to p\pi^-$, we perform the vertex fit by considering both $p\pi^-$ combinations. The $p\pi^-$ combination with the smallest value of $|M(p\pi^-) - m_\Lambda|$, where m_Λ is the Λ nominal mass, is taken as the Λ candidate. The Λ signal region is defined as $|M(p\pi^-) - m_\Lambda| < 0.003$ GeV/c^2. Finally, a vertex fit is performed to the combination of the $\Lambda/\bar{\Lambda}$ and the remaining p for the reactions $\Lambda p \to \Lambda p/\bar{\Lambda} p \to \bar{\Lambda} p$.

To select the signal events of $J/\psi \to \Lambda\bar{\Lambda}$, the invariant mass recoiling against the $\bar{\Lambda}/\Lambda$, $M_{\rm recoil}(\bar{\Lambda}/\Lambda)$, is required to be in the $\Lambda/\bar{\Lambda}$ signal region, defined as $[m_{\Lambda/\bar{\Lambda}} - 0.020, m_{\Lambda/\bar{\Lambda}} + 0.016]$ GeV/c^2, where $M_{\rm recoil}(\bar{\Lambda}/\Lambda) \equiv \sqrt{E_{\rm beam}^2 - |\vec{p}_{\bar{\Lambda}/\Lambda}c|^2}/c^2$, $E_{\rm beam}$ is the $e^\pm$ beam energy, and $\vec{p}_{\bar{\Lambda}/\Lambda}$ is the measured momentum of the $\bar{\Lambda}/\Lambda$ candidate in the e^+e^- rest frame. The main background is $J/\psi \to \Lambda\bar{\Lambda}$, $\Lambda \to p\pi^-$, $\bar{\Lambda} \to \bar{p}\pi^+$, where no scattering of $\Lambda/\bar{\Lambda}$ with a proton from the beam pipe occurred. To suppress this background, the recoil mass of $\bar{\Lambda} p_\Lambda/\Lambda\bar{p}$, $M_{\rm recoil}(\bar{\Lambda} p_\Lambda/\Lambda\bar{p})$, is obtained from the four-momenta of the initial e^+e^- system and the $\bar{\Lambda}/\Lambda$ and $p_\Lambda/\bar{p}$ candidates, where p_Λ is the proton from Λ decays. $M_{\rm recoil}(\bar{\Lambda} p_\Lambda/\Lambda\bar{p})$ should be around the nominal π^-/π^+ mass for this background, so we require $M_{\rm recoil}(\bar{\Lambda} p_\Lambda/\Lambda\bar{p}) < 0$ GeV/c^2 to remove these events. To select those signal events that react with the cooling oil in the beam pipe, the R_{xy} signal region is defined as [3.0, 3.5] cm, taking into account the detector resolution, where R_{xy} is the distance from the reconstructed $\Lambda p/\bar{\Lambda} p$ vertex to the z axis. To remove the events from the reactions between $\Lambda/\bar{\Lambda}$ and ^{197}Au/^{9}Be/^{12}C nuclei, we define the momentum of the proton in the ^{1}H of the cooling oil as $P(p_{\rm oil}) \equiv |\vec{P}_{\Lambda/\bar{\Lambda}} + \vec{P}_p - (\vec{P}_{e^+e^-} - \vec{P}_{\bar{\Lambda}/\Lambda})|$, where $\vec{P}$ represents the momentum of each particle in the lab frame. Because the proton in the ^{1}H of

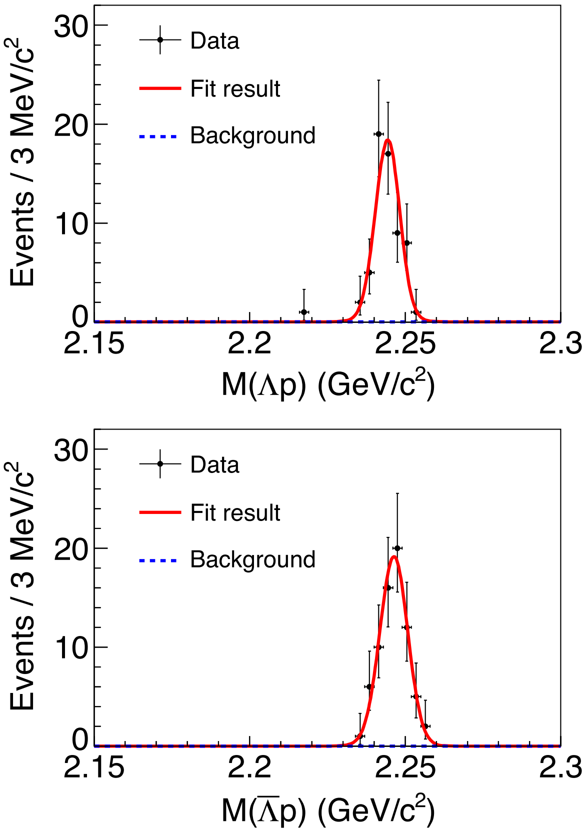

FIG. 2. Distributions of $M(\Lambda p)$ (top) and $M(\bar{\Lambda} p)$ (bottom) of data (black dots with error bars) for the reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$, respectively. The red solid curve is the total fit result and the blue dashed curve is the background component.

TABLE I. Relevant parameters for the differential cross sections, where $\cos\theta_{\Lambda/\bar{\Lambda}}$ is the scattering angle, N_i^{sig} is the number of signal events, ϵ_i is the efficiency, $(d\sigma/d\Omega)$ is the differential cross section, and i represents the different $\cos\theta_{\Lambda/\bar{\Lambda}}$ bins. The first value in parentheses is for $\Lambda p \to \Lambda p$, and the second for $\bar{\Lambda} p \to \bar{\Lambda} p$.

$\cos\theta_{\Lambda/\bar{\Lambda}}$	N_i^{sig}	ϵ_i (%)	$(d\sigma/d\Omega)$ (mb/sr)
$[-0.9, -0.7]$	$(5.0^{+2.6}_{-1.9}, 0.0^{+1.1}_{-0.0})$	(6.94,4.93)	$(1.7^{+0.9}_{-0.7}, 0.0^{+0.5}_{-0.0})$
$(-0.7, -0.5]$	$(1.0^{+1.4}_{-0.7}, 0.0^{+1.1}_{-0.0})$	(14.13,10.44)	$(0.2^{+0.2}_{-0.1}, 0.0^{+0.3}_{-0.0})$
$(-0.5, -0.3]$	$(1.0^{+1.4}_{-0.7}, 1.0^{+1.4}_{-0.7})$	(17.32,13.27)	$(0.2^{+0.2}_{-0.1}, 0.2^{+0.3}_{-0.1})$
$(-0.3, -0.1]$	$(11.0^{+3.7}_{-3.0}, 0.0^{+1.1}_{-0.0})$	(17.74,14.66)	$(1.5^{+0.5}_{-0.4}, 0.0^{+0.2}_{-0.0})$
$(-0.1, 0.1]$	$(6.9^{+3.0}_{-2.3}, 0.0^{+1.1}_{-0.0})$	(19.11,15.79)	$(0.9^{+0.4}_{-0.3}, 0.0^{+0.2}_{-0.0})$
$(0.1,\ 0.3]$	$(5.0^{+2.6}_{-1.9}, 2.0^{+1.8}_{-1.1})$	(19.53,16.82)	$(0.6^{+0.3}_{-0.2}, 0.3^{+0.3}_{-0.2})$
$(0.3,\ 0.5]$	$(12.0^{+3.8}_{-3.1}, 7.0^{+3.0}_{-2.3})$	(19.21,17.68)	$(1.5^{+0.5}_{-0.4}, 1.0^{+0.4}_{-0.3})$
$(0.5,\ 0.7]$	$(13.0^{+3.9}_{-3.3}, 25.0^{+5.3}_{-4.7})$	(19.71,17.60)	$(1.6^{+0.5}_{-0.4}, 3.4^{+0.7}_{-0.6})$
$(0.7,\ 0.9]$	$(6.0^{+2.8}_{-2.1}, 37.0^{+6.4}_{-5.8})$	(9.80,9.93)	$(1.5^{+0.7}_{-0.5}, 9.0^{+1.6}_{-1.4})$

the cooling oil is practically static, while the proton in the ^{197}Au/^{9}Be/^{12}C nuclei has Fermi momentum, the $P(p_{\mathrm{oil}})$ should be around zero for signal processes but hundreds of MeV/c for background processes. To remove these events, the requirement $P(p_{\mathrm{oil}}) < 0.04$ GeV/c is applied.

For the signal reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$ produced from the decay $J/\psi \to \Lambda\bar{\Lambda}$, the center-of-mass energies for the incident $\Lambda/\bar{\Lambda}$ and a static p are all 2.243 GeV/c^2 within a range of ± 0.005 GeV/c^2. Figure 2 shows the $M(\Lambda p)$ and $M(\bar{\Lambda} p)$ distributions from data after the final event selection. Clear enhancements are seen around 2.243 GeV/c^2, corresponding to the reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$, respectively. A detailed study of the J/ψ inclusive MC sample shows that there is no peaking background in the signal region. To determine the signal yield, an unbinned maximum likelihood fit is performed to the $M(\Lambda p)$ distribution and $M(\bar{\Lambda} p)$ distribution, respectively. We use the MC-determined shape convolved with a free Gaussian function to describe the signal, where the yield acts as a free fit parameter. The free Gaussian function is used to describe the difference in the data and signal MC resolutions. The background is described by a uniform distribution with the number of events as free parameter. The fit results are shown in Fig. 2. The signal yields returned by the fits are $N^{\mathrm{sig}}_{\Lambda p} = 60.9 \pm 7.8$ and $N^{\mathrm{sig}}_{\bar{\Lambda} p} = 72.0 \pm 8.5$ for the reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$, respectively, and the goodness of the fits for the two reactions are $\chi^2/\mathrm{ndf} = 4.8/4 = 1.2$ and $0.8/4 = 0.2$ without considering empty bins.

To extract the differential cross sections for the two reactions, we need the signal yields as a function of $\cos\theta_{\Lambda/\bar{\Lambda}}$, where $\theta_{\Lambda/\bar{\Lambda}}$ is the scattering angle of the scattered $\Lambda/\bar{\Lambda}$ in the $\Lambda p/\bar{\Lambda} p$ rest frames with the z axis defined by the incident $\Lambda/\bar{\Lambda}$ momentum. Because the efficiency is very low and it is hard to obtain accurate experimental information near the regions $\cos\theta_{\Lambda/\bar{\Lambda}} = \pm 1$ due to the low momentum of scattered $\Lambda/\bar{\Lambda}$ or p, the measurements are restricted to $-0.9 \le \cos\theta_{\Lambda/\bar{\Lambda}} \le 0.9$. To obtain the number of signal events, we perform a simultaneous fit to the $M(\Lambda p)$ and $M(\bar{\Lambda} p)$ distributions in nine different $\cos\theta_{\Lambda/\bar{\Lambda}}$ regions, where the signal shape and background shape are the same as mentioned above. The obtained number of signal events in the nine $\cos\theta_{\Lambda/\bar{\Lambda}}$ regions are summarized in Table I. It is worth mentioning that no events survived in the $-1.0 < \cos\theta_{\Lambda/\bar{\Lambda}} < -0.9$ and the $0.9 < \cos\theta_{\Lambda/\bar{\Lambda}} < 1.0$ regions for data.

Using the same method as in Ref. [23], the cross sections of the reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$ can be determined, the only difference is that we use the proton in the ^{1}H of the cooling oil of the beam pipe as the target material. The total elastic cross sections are calculated with

$$\sigma(\Lambda p \to \Lambda p/\bar{\Lambda} p \to \bar{\Lambda} p) = \frac{N^{\mathrm{sig}}_{\Lambda p/\bar{\Lambda} p}}{\epsilon_{\Lambda p/\bar{\Lambda} p} \mathcal{B} \mathcal{L}_{\mathrm{eff}}}, \qquad (1)$$

where $\epsilon_{\Lambda p/\bar{\Lambda} p} = [\sum_i \epsilon_i (d\sigma/d\Omega)_i / \sum_i (d\sigma/d\Omega)_i]$ is the weighted selection efficiency according to the differential

TABLE II. Input parameters for the cross section calculations. The first value in brackets is for $\Lambda p \to \Lambda p$, and the second is for $\bar{\Lambda} p \to \bar{\Lambda} p$.

Parameter	Result
$N^{\mathrm{sig}}_{\Lambda p/\bar{\Lambda} p}$	$(60.9 \pm 7.8, 72.0 \pm 8.5)$
$\epsilon_{\Lambda p/\bar{\Lambda} p}$	$(15.29\%, 12.55\%)$
$\mathcal{B}$	$(40.8321 \pm 0.4518)\%$ [1]
$N_{J/\psi}$	$(10.087 \pm 0.044) \times 10^9$ [21]
$\mathcal{B}_{J/\psi}$	$(0.189 \pm 0.009)\%$ [1]
α	0.475 ± 0.004 [30]
L	(7.89 ± 0.06) cm [1]
E_{beam}	1.5485 GeV
$m_{\Lambda/\bar{\Lambda}}$	(1.115683 ± 0.000006) GeV$/c^2$ [1]
a	3.23 cm [20]
b	3.31 cm [20]
N_H	7.35×10^{22} cm^{-3}

cross section distribution, which will be introduced later. $\mathcal{B}$ is the product of the branching ratios of the intermediate states, defined as $\mathcal{B} \equiv \mathcal{B}(\Lambda \to p\pi^-)\mathcal{B}(\bar{\Lambda} \to \bar{p}\pi^+)$, and $\mathcal{L}_{\mathrm{eff}}$ is the effective luminosity of the reaction of the $\Lambda/\bar{\Lambda}$ flux produced from $J/\psi \to \Lambda\bar{\Lambda}$ with the target material:

$$\mathcal{L}_{\mathrm{eff}} = \frac{N_{J/\psi}\mathcal{B}_{J/\psi}}{2+\frac{2}{3}\alpha}\int_a^b \int_0^\pi (1+\alpha\cos^2\theta)e^{-\frac{x}{\sin\theta\beta\gamma L}} N_H \mathrm{d}\theta \mathrm{d}x. \quad (2)$$

In the integral of this formula, the angular distribution of the $\Lambda/\bar{\Lambda}$ flux, the attenuation of the $\Lambda/\bar{\Lambda}$ flux, and the number of target nuclei are considered. $N_{J/\psi}$ is the number of J/ψ events, $\mathcal{B}_{J/\psi}$ is the branching fraction of $J/\psi \to \Lambda\bar{\Lambda}$, and α is the parameter of the angular distribution of $J/\psi \to \Lambda\bar{\Lambda}$, $\beta\gamma \equiv (\sqrt{E^2_{\mathrm{beam}} - m^2_{\Lambda/\bar{\Lambda}}c^4}/m_{\Lambda/\bar{\Lambda}}c^2)$ is the ratio of the momentum to the mass of the $\Lambda/\bar{\Lambda}$, and $L \equiv c\tau$ is the product of the speed of light and the mean lifetime of the $\Lambda/\bar{\Lambda}$ [1]. N_H is the number of target nuclei per unit volume, a and b are the distances from the inner surface and outer surface of the cooling oil in the beam pipe to the z axis, θ and x are the angle and distance to the z axis, as shown in Fig. 1. The beam pipe can be regarded as infinitely long with respect to the product of $\beta\gamma L$ for $\Lambda/\bar{\Lambda}$. The parameters are listed in Table II, and the corresponding total elastic cross sections in $-0.9 \le \cos\theta_{\Lambda/\bar{\Lambda}} \le 0.9$ are measured to be $\sigma(\Lambda p \to \Lambda p) = (12.2 \pm 1.6_{\mathrm{stat}} \pm 1.1_{\mathrm{syst}})$ and $\sigma(\bar{\Lambda} p \to \bar{\Lambda} p) = (17.5 \pm 2.1_{\mathrm{stat}} \pm 1.6_{\mathrm{syst}})$ mb at a $\Lambda/\bar{\Lambda}$ momentum of 1.074 GeV$/c$ within a range of ± 0.017 GeV$/c$.

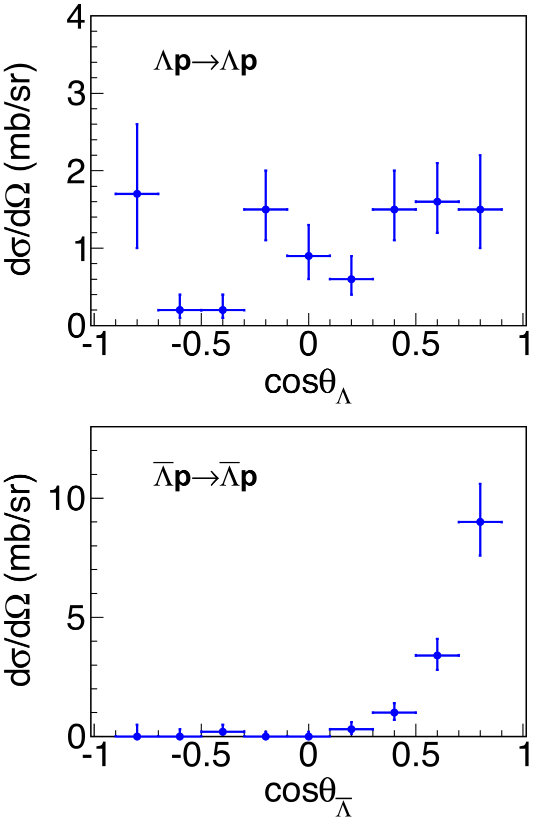

FIG. 3. Differential cross sections of the reactions $\Lambda p \to \Lambda p$ (top) and $\bar{\Lambda} p \to \bar{\Lambda} p$ (bottom) at the $\Lambda/\bar{\Lambda}$ momentum of around 1.074 GeV$/c$.

The differential cross sections for the reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$ are calculated with

$$\left(\frac{d\sigma}{d\Omega}\right)_i = \frac{N_i^{\mathrm{sig}}}{\epsilon_i \mathcal{B} \mathcal{L}_{\mathrm{eff}} \Delta\Omega}, \quad (3)$$

where N_i^{sig} and ϵ_i are the number of signal events and efficiency, i represents different $\cos\theta_{\Lambda/\bar{\Lambda}}$ bins, and $\Delta\Omega = 2\pi\Delta\cos\theta_{\Lambda/\bar{\Lambda}} = 0.4\pi$ represents the solid angle. The measured results are listed in Table I and shown in Fig. 3. We can see there is a slight tendency of forward scattering for $\Lambda p \to \Lambda p$, while there is a strong forward peak for $\bar{\Lambda} p \to \bar{\Lambda} p$. The different behaviors indicate that the reaction mechanisms of these two processes are different.

We also tested an extrapolation for the regions of $|\cos\theta_{\Lambda/\bar{\Lambda}}| > 0.9$ for the differential cross sections of $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$ to determine the total elastic cross sections. For the reaction $\Lambda p \to \Lambda p$, we assume the differential cross sections in $-1.0 < \cos\theta_\Lambda < -0.9$ and $0.9 < \cos\theta_\Lambda < 1.0$ to be the same as those in neighboring bins. For the reaction $\bar{\Lambda} p \to \bar{\Lambda} p$, the differential cross section is fitted using a piecewise polynomial function, which is a constant for $\cos\theta_{\bar{\Lambda}} \le 0$ and a third-order polynomial function for $\cos\theta_{\bar{\Lambda}} \ge 0$. The differential cross section in the regions of $|\cos\theta_{\bar{\Lambda}}| > 0.9$ is obtained according to the fit function. Therefore, the total elastic cross

TABLE III. Summary of systematic uncertainties (in %).

Source	$\sigma(\Lambda p \to \Lambda p/\bar{\Lambda} p \to \bar{\Lambda} p)$
Tracking efficiency	5.0
PID efficiency	5.0
Track number	2.2
Branching fractions	4.9
e^+e^- interaction point	2.0
Sum	9.1

sections integrated over the full angular region are determined to be $\sigma_t(\Lambda p \to \Lambda p) = (14.2 \pm 1.8_{\text{stat}} \pm 1.3_{\text{syst}})$ and $\sigma_t(\bar{\Lambda} p \to \bar{\Lambda} p) = (27.4 \pm 3.2_{\text{stat}} \pm 2.5_{\text{syst}})$ mb. The result of the total elastic cross section on the reaction $\Lambda p \to \Lambda p$ is consistent with those measured from other experiments [10–12,31–42]. The strong forward rise of the differential cross section of $\bar{\Lambda} p \to \bar{\Lambda} p$ is compatible with the expectation for the case of scattering in the presence of a strong absorption [43–45], which is given by the annihilation part of the potential. Especially, this behavior is very similar to $\bar{p}p$ elastic scattering in a comparable incident momentum region [44], in contrast, such a strong forward rise does not appear in pp elastic scattering [46]. This indicates that the strong absorption mechanism is not only important in $\bar{\text{N}}$N scattering, but also in $\bar{\text{Y}}$N scattering. If we assume the reaction $\bar{\Lambda} p \to \bar{\Lambda} p$ is a pure "black sphere" scattering, the total elastic cross section is given by $\sigma_t(\bar{\Lambda} p \to \bar{\Lambda} p) = \pi R^2$ [43], where R is the interaction radius. This gives $R = (0.93 \pm 0.07)$ fm, which is comparable to the proton radius [1].

The sources of systematic uncertainties related to the measured cross sections are discussed in the following. The uncertainties in the tracking efficiency and PID efficiency are 1% per track [23]. The uncertainty of the track number requirement is estimated with the control sample $J/\psi \to \Lambda\bar{\Lambda} \to p\pi^- \bar{p}\pi^+$. The uncertainties for the branching fractions are taken from the PDG [1]. To estimate the uncertainty from the position of the e^+e^- interaction point, we change the integral range by ±0.1 cm, which is from (a, b) to $(a + 0.1, b + 0.1)$ or $(a - 0.1, b - 0.1)$, and the larger difference in the result is taken as the uncertainty. The systematic uncertainties from $\Lambda/\bar{\Lambda}$ mass windows, $M_{\text{recoil}}(\bar{\Lambda} p_\Lambda)/M_{\text{recoil}}(\Lambda\bar{p})$ requirement, R_{xy} requirement and $P(p_{\text{oil}})$ requirement are tested using a Barlow test method [24], and these items can be considered negligible. The systematic uncertainties from the fit procedure, the number of J/ψ events, the angular distribution of $J/\psi \to \Lambda\bar{\Lambda}$, and the Λ mean lifetime are all less than 1% and can be ignored. A summary of the main systematic uncertainties is presented in Table III, and the total systematic uncertainty is obtained by adding all the individual components in quadrature.

In summary, using $(10.087 \pm 0.044) \times 10^9$ J/ψ events collected with the BESIII detector operating at the BEPCII storage ring, the reactions $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$ are measured, where $\Lambda/\bar{\Lambda}$ are from the process $J/\psi \to \Lambda\bar{\Lambda}$ and p is from the cooling oil in the beam pipe. The cross sections in $-0.9 \leq \cos\theta_{\Lambda/\bar{\Lambda}} \leq 0.9$ are measured to be $\sigma(\Lambda p \to \Lambda p) = (12.2 \pm 1.6_{\text{stat}} \pm 1.1_{\text{syst}})$ and $\sigma(\bar{\Lambda} p \to \bar{\Lambda} p) = (17.5 \pm 2.1_{\text{stat}} \pm 1.6_{\text{syst}})$ mb at the $\Lambda/\bar{\Lambda}$ momentum of 1.074 GeV/c within a range of ±0.017 GeV/c. Furthermore, the differential cross sections of the two reactions are also measured. There is a slight tendency of forward scattering for $\Lambda p \to \Lambda p$, while a strong forward peak for $\bar{\Lambda} p \to \bar{\Lambda} p$ is observed. If we make an extrapolation for the regions of $|\cos\theta_{\Lambda/\bar{\Lambda}}| > 0.9$ for the differential cross sections of $\Lambda p \to \Lambda p$ and $\bar{\Lambda} p \to \bar{\Lambda} p$, the total elastic cross sections integrated over the full angular region are determined to be $\sigma_t(\Lambda p \to \Lambda p) = (14.2 \pm 1.8_{\text{stat}} \pm 1.3_{\text{syst}})$ and $\sigma_t(\bar{\Lambda} p \to \bar{\Lambda} p) = (27.4 \pm 3.2_{\text{stat}} \pm 2.5_{\text{syst}})$ mb. These constitute the first result of $\bar{\text{Y}}$N scattering, and will serve as input for the theoretical understanding of the (anti) hyperon-nucleon interaction. This work is the first study of (anti)hyperon-nucleon elastic scattering at an electron-positron collider, and demonstrates the feasibility for studying other antihyperons, such as $\bar{\Sigma} p \to \bar{\Sigma} p$ and $\bar{\Xi} p \to \bar{\Xi} p$. The momentum dependence of these cross sections could be studied at a future super tau-charm factory [47,48] by exploiting multibody processes or other charmonium decays.

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