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## Growth and properties of $\text{Cu}_3\text{N}$ films and $\text{Cu}_3\text{N}/\gamma\text{-Fe}_4\text{N}$ bilayers

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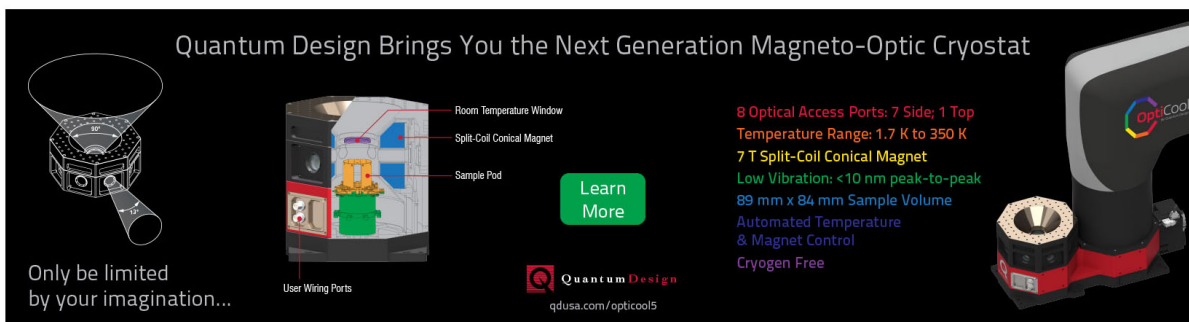
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## Growth and properties of $\text{Cu}_3\text{N}$ films and $\text{Cu}_3\text{N}/\gamma\text{-Fe}_4\text{N}$ bilayers

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Copper nitride films were grown by molecular-beam epitaxy of copper in the presence of nitrogen from a radio-frequency atomic source on (001)  $\gamma\text{-Fe}_4\text{N}/(001)\text{MgO}$  or directly on MgO substrates. The structural properties of the  $\text{Cu}_3\text{N}$  films were found to be very dependent on the substrate and on the deposition temperature. At optimal growth conditions, the  $\text{Cu}_3\text{N}$  films grow epitaxial on both substrates. The  $\text{Cu}_3\text{N}$  films grown on MgO were characterized optically to be insulators with an energy gap of 1.65 eV. On  $\gamma\text{-Fe}_4\text{N}$ ,  $\text{Cu}_3\text{N}$  films with a thickness of only 6 nm, were grown as closed layers, epitaxial and rather smooth (root-mean-square roughness of 0.7 nm). This material has ideal properties to be used as a barrier in low resistance magnetic tunnel junctions. © 2002 American Institute of Physics. [DOI: 10.1063/1.1459116]

Transition metal nitrides show a wide variety of properties and applications. Among them, copper nitrides have attracted considerable attention as a new material for optical storage devices, based on the rather low decomposition temperature.<sup>1-3</sup>

According to theoretical calculations,<sup>4</sup>  $\text{Cu}_3\text{N}$  is an insulator, with an energy gap of  $\sim 0.9$  eV, whereas experimental values ranging from 0.8 to 1.9 eV were found. The structure of  $\text{Cu}_3\text{N}$  is, in itself, rather interesting. With the cubic anti- $\text{ReO}_3$  structure, the face-centered-cubic close-packed sites are vacant and, therefore, it is possible to insert another metallic atom (e.g., Cu and Pd). With the insertion of Cu atoms, the  $\text{Cu}_3\text{N}$  lattice undergoes a small ( $\sim 1.13\%$ ) expansion, and the  $\text{Cu}_4\text{N}$  phase is formed.<sup>5,6</sup> Contrary to  $\text{Cu}_3\text{N}$ , this compound is believed to be a conductor. Similarly,  $\text{Cu}_3\text{NPd}$  compounds were synthesized, also with metallic character.<sup>4,7</sup>

Up until now, films of  $\text{Cu}_3\text{N}$  or  $\text{Cu}_3\text{N}/\text{Cu}$  were reported to be grown by different methods, such as rf magnetron sputtering,<sup>1,3,8-12</sup> dc sputtering,<sup>13</sup> or ion-assisted deposition.<sup>2</sup> The obtained films were polycrystalline on substrates like Pt/MgO,  $\text{Al}_2\text{O}_3$ , Si, or glass and amorphous on MgO and  $\text{SrTiO}_3$ .<sup>8</sup> The properties of the films were found to be very dependent on substrate and growth conditions.

In this letter, the epitaxial growth and properties of thin films of  $\text{Cu}_3\text{N}$  on  $\gamma\text{-Fe}_4\text{N}/\text{MgO}$  and MgO substrates are reported. This material, with a rather low energy gap, might be used as a barrier material in low resistance magnetic tunnel junctions. Additionally, the excellent match between  $\text{Cu}_3\text{N}$ , an insulator, and  $\gamma\text{-Fe}_4\text{N}$ , a ferromagnetic conductor, opens the perspective of developing an all-nitride, all-epitaxial magnetic tunnel junction.

Thin films of copper nitride were grown in an ultrahigh vacuum (UHV) chamber (base pressure of  $10^{-10}$  mbar) by molecular-beam epitaxy (MBE) of copper in the presence of atomic nitrogen obtained from a rf atomic source. As substrates, we used (001)  $\gamma\text{-Fe}_4\text{N}$  thin films grown on (001) MgO. The iron nitride films were grown in the same chamber, in a similar fashion as the copper nitride films. In previ-

ous work, we proved that this method is very effective to grow epitaxial iron nitride films.<sup>14,15</sup> As reported earlier, the optimal substrate temperature for the growth of  $\gamma\text{-Fe}_4\text{N}$  films on (001) MgO substrates was 400 °C. At such substrate temperatures, the obtained films were of high crystal quality and rather smooth [for a  $\sim 33$  nm film the root-mean-square (rms) roughness was 0.4 nm]. Based on the very good lattice match, the  $\gamma\text{-Fe}_4\text{N}$  phase (cubic, 3.795 Å) is an ideal substrate for the epitaxial growth of  $\text{Cu}_3\text{N}$  (cubic, 3.817 Å).

Additionally,  $\text{Cu}_3\text{N}$  films were also grown directly on (001) MgO substrates, in a similar way. Here the lattice mismatch is  $\sim 10\%$ . In all cases, before growth, the MgO substrates were cleaned by annealing at 600 °C in  $10^{-6}$  mbar  $\text{O}_2$ . The copper was evaporated from a Knudsen cell, at rather low deposition rates (70–130 s/Å) on  $\gamma\text{-Fe}_4\text{N}/\text{MgO}$  substrates heated at 150 °C and on MgO substrates heated at 150 °C or 250 °C. The rf atomic source was operated with pure nitrogen at a gas pressure of  $5 \times 10^{-2}$  mbar and also with mixtures of nitrogen and hydrogen. The rf power applied was 60 W. The pressure in the UHV chamber during growth was  $10^{-7}$  mbar.

The crystal structure was investigated by means of x-ray diffraction in a standard  $\theta$ - $2\theta$  geometry (XRD) and also by performing texture measurements in a Philips X'Pert MRD system [two-dimensional-XRD (2D-XRD)]. Both measurements were done using the Cu  $K\alpha$  radiation ( $\lambda = 1.54$  nm). The surface morphology was probed with a NanoScope IIIa atomic force microscope (AFM). The optical properties were obtained using a Woollam variable angle spectroscopy ellipsometer. The thickness of the films was measured by Ruthenford backscattering spectroscopy.

With only nitrogen in the rf atomic source, the optimal deposition temperature to obtain epitaxial layers was found to be 150 °C on  $\gamma\text{-Fe}_4\text{N}$  substrates and 150 °C–250 °C on MgO. The upper limit of the deposition temperature is dictated by the N uptake in the sample: at too high temperatures no nitrogen is taken up and a pure Cu layer is grown. Also, the formation of  $\text{Cu}_3\text{N}$  films was found not to be critically dependent on the growth rate. Contrary to the growth of

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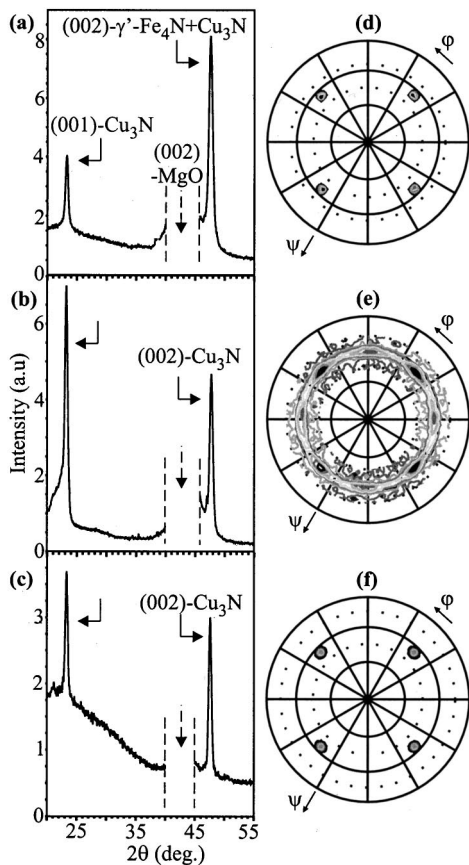


FIG. 1. X-Ray  $\theta$ - $2\theta$  scans: (a) 17 nm  $\text{Cu}_3\text{N}/15$  nm  $\gamma'$ - $\text{Fe}_4\text{N}/\text{MgO}$  grown at  $150^\circ\text{C}$ , (b) 18.6 nm  $\text{Cu}_3\text{N}/\text{MgO}$  grown at  $150^\circ\text{C}$ , (c) 15.7 nm  $\text{Cu}_3\text{N}/\text{MgO}$  grown at  $250^\circ\text{C}$ . The corresponding pole  $\psi$ - $\phi$  scans are shown in (d), (e), and (f). In all pole scans, the  $2\theta$  angle is fixed for the (111) reflection of  $\text{Cu}_3\text{N}$ .

$\gamma'$ - $\text{Fe}_4\text{N}$ ,<sup>15</sup> we found that the formation of pure  $\text{Cu}_3\text{N}$  films is less effective when hydrogen is present in a mixture with nitrogen in the rf atomic source, as compared with the case when only nitrogen is used. Possibly, part of the needed nitrogen for the formation of  $\text{Cu}_3\text{N}$  combines with hydrogen, and finally desorbs as  $\text{NH}_3$ .

Typical results of the XRD measurements for different  $\text{Cu}_3\text{N}$  samples are shown in Fig. 1, as follows: Fig. 1(a) for the bilayer structure, 17 nm  $\text{Cu}_3\text{N}/15$  nm  $\gamma'$ - $\text{Fe}_4\text{N}/\text{MgO}$  with the copper nitride film grown at  $150^\circ\text{C}$ , Fig. 1(b) for a 18.6 nm  $\text{Cu}_3\text{N}$  film grown at  $150^\circ\text{C}$  on  $\text{MgO}$  and (c) for a 15.7 nm  $\text{Cu}_3\text{N}$  film grown at  $250^\circ\text{C}$  also on  $\text{MgO}$ . When  $\text{Cu}_3\text{N}$  is grown on a  $\gamma'$ - $\text{Fe}_4\text{N}$  underlayer, the (002) reflections overlap, which results in a slightly broader peak at  $2\theta = 47.76^\circ$ . However,  $\text{Cu}_3\text{N}$  can be recognized by the presence of a rather intense (001) peak at  $2\theta = 23.3^\circ$ , as compared with the one of  $\gamma'$ - $\text{Fe}_4\text{N}$  (for  $\text{Cu}_3\text{N}$  the intensity ratio of the (001) and (002) reflections is 1.16 whereas for  $\gamma'$ - $\text{Fe}_4\text{N}$  this ratio is only 0.028). With bare  $\text{MgO}$  substrates, the (002) and (001) reflections of  $\text{Cu}_3\text{N}$  are clearly visible in  $\theta$ - $2\theta$  scans. In all scans, only (00 $k$ ) reflections corresponding to copper nitride and iron nitride are present. Therefore, we conclude that all copper nitride films have a well-defined orientation with the [100] direction normal to the (001) substrate.

To determine the in-plane orientation of the film with respect to the crystal orientation of the substrate, we performed 2D-XRD  $\psi$ - $\phi$  pole scans with  $2\theta$  fixed at the (111)

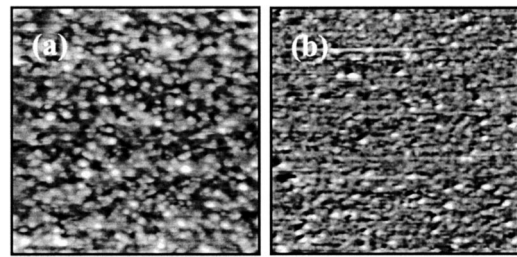


FIG. 2. AFM topographical images of  $\text{Cu}_3\text{N}$  films of different thickness ( $t$ ) grown under similar conditions on  $\gamma'$ - $\text{Fe}_4\text{N}$  substrates. (a)  $t = 17$  nm,  $1 \times 1 \mu\text{m}^2$  image (vertical range 25 nm; rms roughness  $\sim 2.6$  nm). (b)  $t = 6$  nm,  $1 \times 1 \mu\text{m}^2$  image (vertical range 10 nm; rms roughness  $\sim 0.7$  nm).

reflection of  $\text{Cu}_3\text{N}$ . Note that also the  $\text{Cu}_3\text{N}$  (111) reflection is very close to the (111) reflection of  $\gamma'$ - $\text{Fe}_4\text{N}$ . The results corresponding to the three samples are presented in Figs. 1(d)–1(f). As shown in Fig. 1(d), only four sharp peaks corresponding to the (111) reflections of both  $\text{Cu}_3\text{N}$  and  $\gamma'$ - $\text{Fe}_4\text{N}$  are present. Therefore, the copper nitride film grown at a deposition temperature of only  $150^\circ\text{C}$ , on a matching substrate, is epitaxial and monocrystalline with the same symmetry as the iron nitride layer with respect to the  $\text{MgO}$  substrate (no in-plane rotation). For the films grown on  $\text{MgO}$  substrates, the corresponding pole scans are shown in Figs. 1(e) and 1(f). In this case, at the same deposition temperature of  $150^\circ\text{C}$  as for  $\text{Cu}_3\text{N}$  films on  $\gamma'$ - $\text{Fe}_4\text{N}$  (but on a different substrate), a more complex pattern with eight (111) intensity maxima is present [Fig. 1(e)]. The intensity in between the maxima is due to background. This pattern corresponds to an epitaxial film, with two preferred in-plane orientations, instead of only one. At higher deposition temperatures, ( $250^\circ\text{C}$ ), the  $\text{Cu}_3\text{N}$  films are epitaxial and monocrystalline [Fig. 1(f)]. The difference in structure of the  $\text{Cu}_3\text{N}$  films on  $\text{MgO}$  obtained at different deposition temperatures could be explained on the basis of the high mismatch present between the film and the substrate ( $\sim 10\%$ ) combined with a higher mobility of the Cu atoms at a higher deposition temperature. For the  $\text{Cu}_3\text{N}$  films grown on a  $\text{MgO}$  substrate we found from the  $2\theta$  values corresponding to the (002) and (111) reflections of  $\text{Cu}_3\text{N}$ , that a 18.6 nm thick film grown at  $150^\circ\text{C}$ , is fully relaxed, whereas the 15.7 nm thick film grown at  $250^\circ\text{C}$ , shows a slight tetragonal distortion, with a small in-plane contraction and an out-of-plane expansion.

The surface morphology of  $\text{Cu}_3\text{N}$  films of different thickness (6 nm and 17 nm) grown on  $\gamma'$ - $\text{Fe}_4\text{N}$  substrates was studied in air with AFM in contact mode. For a thickness of  $\sim 17$  nm, the  $\text{Cu}_3\text{N}$  films were found to be rough on a nm scale [Fig. 2(a)]. However, for the  $\text{Cu}_3\text{N}$  film of 6 nm thickness, the surface was rather smooth. [Fig. 2(b)]. Despite the rather low deposition temperature, the surface of the film has a uniform granular appearance with a rms roughness of only 0.7 nm. The difference in smoothness between the thin (6 nm) and the thicker (17 nm)  $\text{Cu}_3\text{N}$  film could be explained assuming the Stranski–Krastanov growth mode, with a critical thickness for the transition from the 2D growth regime (layer by layer) to the three-dimensional (3D) growth regime (islands) between 6 and 17 nm. Similar values (and higher) of the critical thickness for the transition from 2D to 3D growth were observed also in different systems (eg.



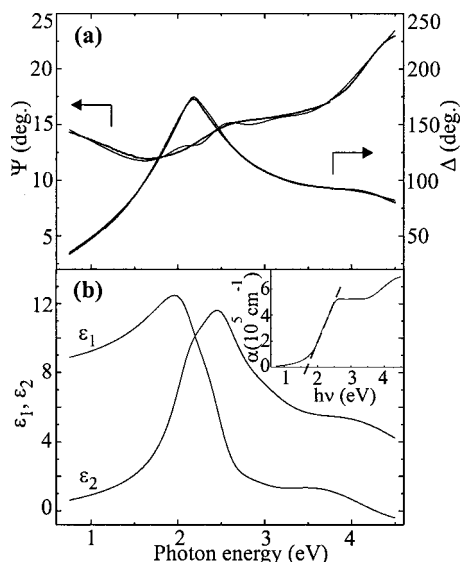


FIG. 3. (a) Measured ellipsometric parameters  $\Psi$  and  $\Delta$  (thick dotted line) together with the fit of the experimental data (thin solid line) for a  $\text{Cu}_3\text{N}$  thin film ( $\sim 18.6$  nm) grown at  $150^\circ\text{C}$  on a MgO substrate. (b) The real ( $\epsilon_1$ ) and imaginary ( $\epsilon_2$ ) parts of the dielectric function. The inset shows the absorption coefficient ( $\alpha$ ) from which the optical energy gap is estimated to be  $\sim 1.65$  eV.

$\text{YBa}_2\text{Cu}_3\text{O}_{7-\delta}$ ). Additional x-ray photoelectron spectroscopy measurements showed, as expected, that the  $\text{Cu}_3\text{N}$  films grow as complete layers for such thin films.

Optical measurements were performed at room temperature in the 0.8–4.5 eV energy range. All spectra were taken at an angle of incidence of  $70^\circ$  and a fixed polarization angle of  $45^\circ$ . Figure 3(a) shows the measured ellipsometric parameters  $\psi$  and  $\Delta$  for a  $\text{Cu}_3\text{N}$  thin film ( $\sim 18.6$  nm) grown at  $150^\circ\text{C}$  on MgO. The experimental data were fitted taking into account the thickness of the  $\text{Cu}_3\text{N}$  film and the optical properties of the MgO substrate, with Lorentzian dielectric functions. The results of the fit are shown together with the experimental data. From the fit, the real ( $\epsilon_1$ ) and imaginary ( $\epsilon_2$ ) parts of the dielectric function were calculated<sup>16</sup> [Fig. 3(b)]. The dependence of  $\epsilon_1$  and  $\epsilon_2$  on energy, typical for an insulator, agrees well with previously reported data. However, in our case, the magnitude of  $\epsilon_1$  and  $\epsilon_2$  is slightly higher than the values reported (all for much thicker films). This might be an effect correlated with the thickness of the films, which is much lower in our case. From the dependence of the absorption coefficient on energy as derived from our data [inset Fig. 3(b)] we estimate an energy gap of 1.65 eV. This value is in the range of data previously measured for thicker  $\text{Cu}_3\text{N}$  films grown on different substrates. For the  $\text{Cu}_3\text{N}$  film grown on MgO at  $250^\circ\text{C}$ , similar results were obtained, although here a small absorption peak was present below the optical gap. Such an effect, also seen in other materials,<sup>17</sup> and known as the Urbach tail, might be due to absorption by impurities (in our case Cu).

Additional resistivity measurements confirmed the insulating character of all the  $\text{Cu}_3\text{N}$  films. Values in the order of  $\sim 1000 \Omega \text{ cm}$  were measured. Due to the complexity of a multilayer system for optical spectroscopy measurements and also because of lack of optical data for  $\gamma'$ - $\text{Fe}_4\text{N}$  films, we could only characterize optically the  $\text{Cu}_3\text{N}$  films grown directly on MgO substrates.

In summary, this letter reports on the epitaxial growth of  $\text{Cu}_3\text{N}/\gamma'$ - $\text{Fe}_4\text{N}$  bilayers and  $\text{Cu}_3\text{N}$  layers on (001) MgO substrates by MBE of copper in the presence of nitrogen from a rf atomic source. The crystal structure of the films was found to be dependent on the deposition temperature, as well as on the substrate. Films of  $\text{Cu}_3\text{N}$  grow as closed epitaxial and monocrystalline layers on epitaxial  $\gamma'$ - $\text{Fe}_4\text{N}$ . Moreover,  $\text{Cu}_3\text{N}$  films with a thickness in the range of interest for tunnel junction applications were found to be rather smooth (rms roughness  $\sim 0.7$  nm). On MgO substrates, the films are also epitaxial. If higher deposition temperatures ( $250^\circ\text{C}$ ) are used, the films are also monocrystalline.

The  $\text{Cu}_3\text{N}$  films grown on MgO substrates were characterized optically to be insulators, with an optical energy gap of 1.65 eV. The present growth technique proved its potential in growing high quality insulating films, which might be good candidates for barriers in magnetic tunnel junctions.

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